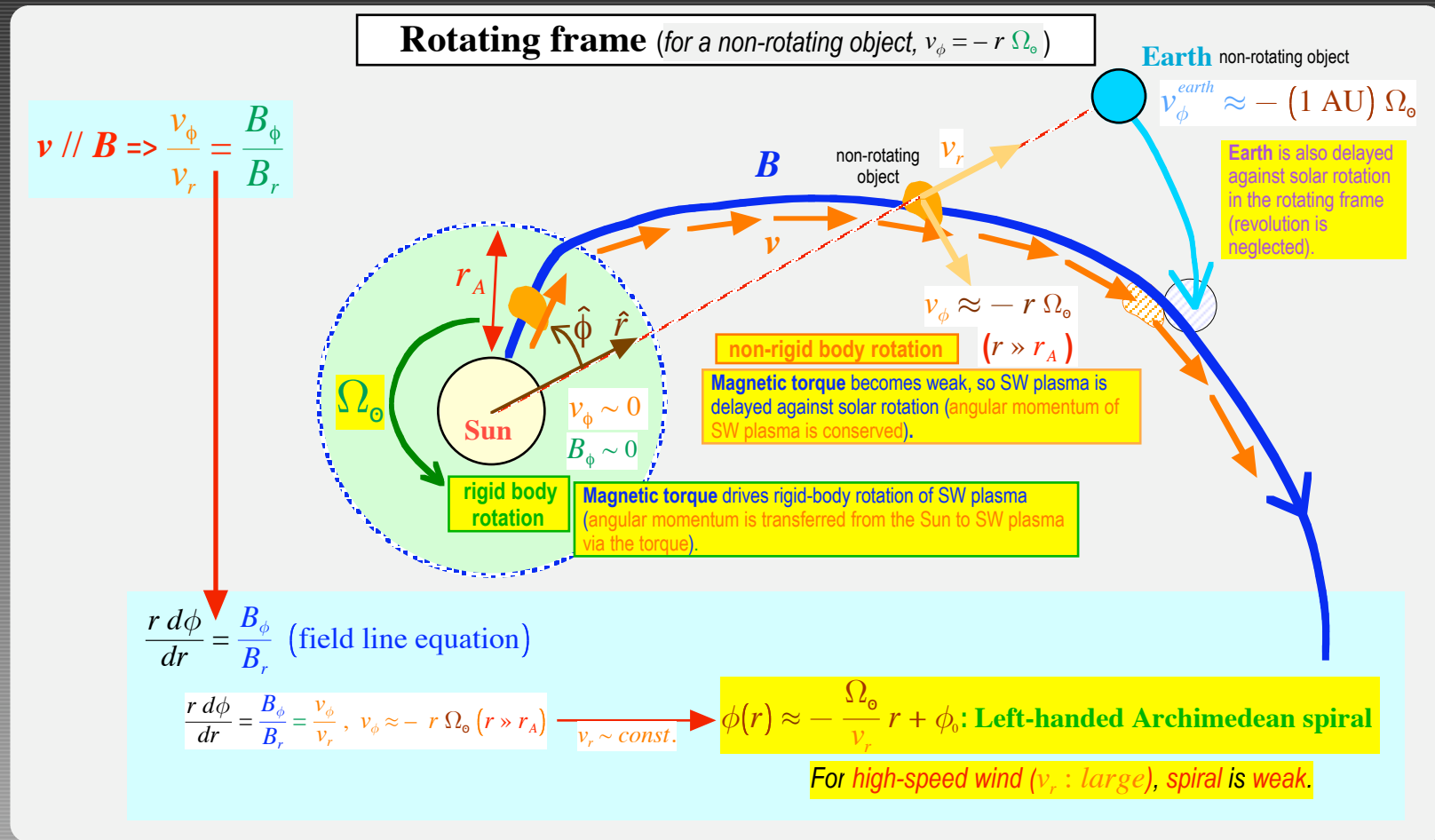


# Archimedean spiral structure of IMF: $\phi(r) \propto r$

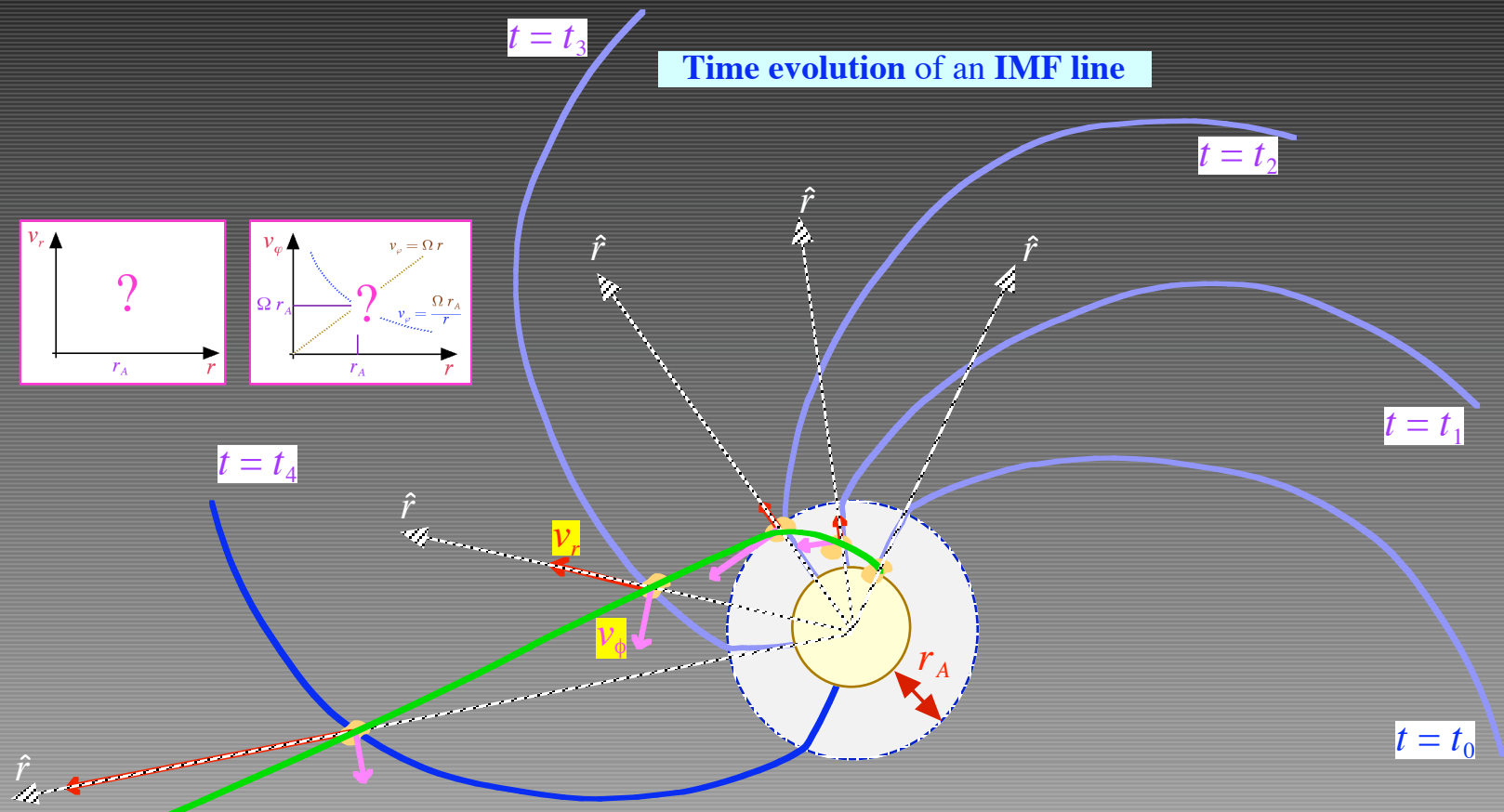
We assume a static IMF frozen into a SW plasma in a frame rotating at the solar rotation rate  $\Omega_\odot$ . In this frame, flow velocity of the plasma is parallel to the IMF:

$$\mathbf{v} \parallel \mathbf{B} \Rightarrow \frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B}) = \mathbf{0}$$



In inertial frame...

Time evolution of an IMF line

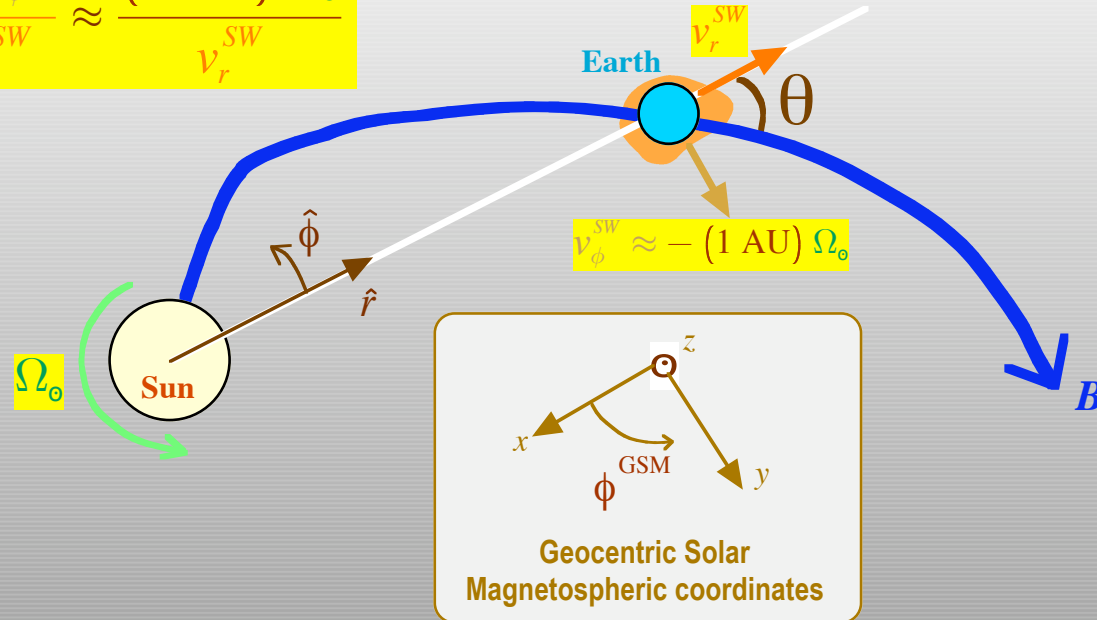


Trajectory of a SW plasma blob

For  $R_0 < r \lesssim r_A$  ... rigid-body rotation  
For  $r_A < r$  ... angular momentum conservation

## Spiral angle of IMF near an Earth orbit

$$\tan \theta \equiv \frac{-B_{\phi}^{IMF}}{B_r^{IMF}} = \frac{-v_{\phi}^{SW}}{v_r^{SW}} \approx \frac{(1 \text{ AU}) \Omega_{\odot}}{v_r^{SW}}$$

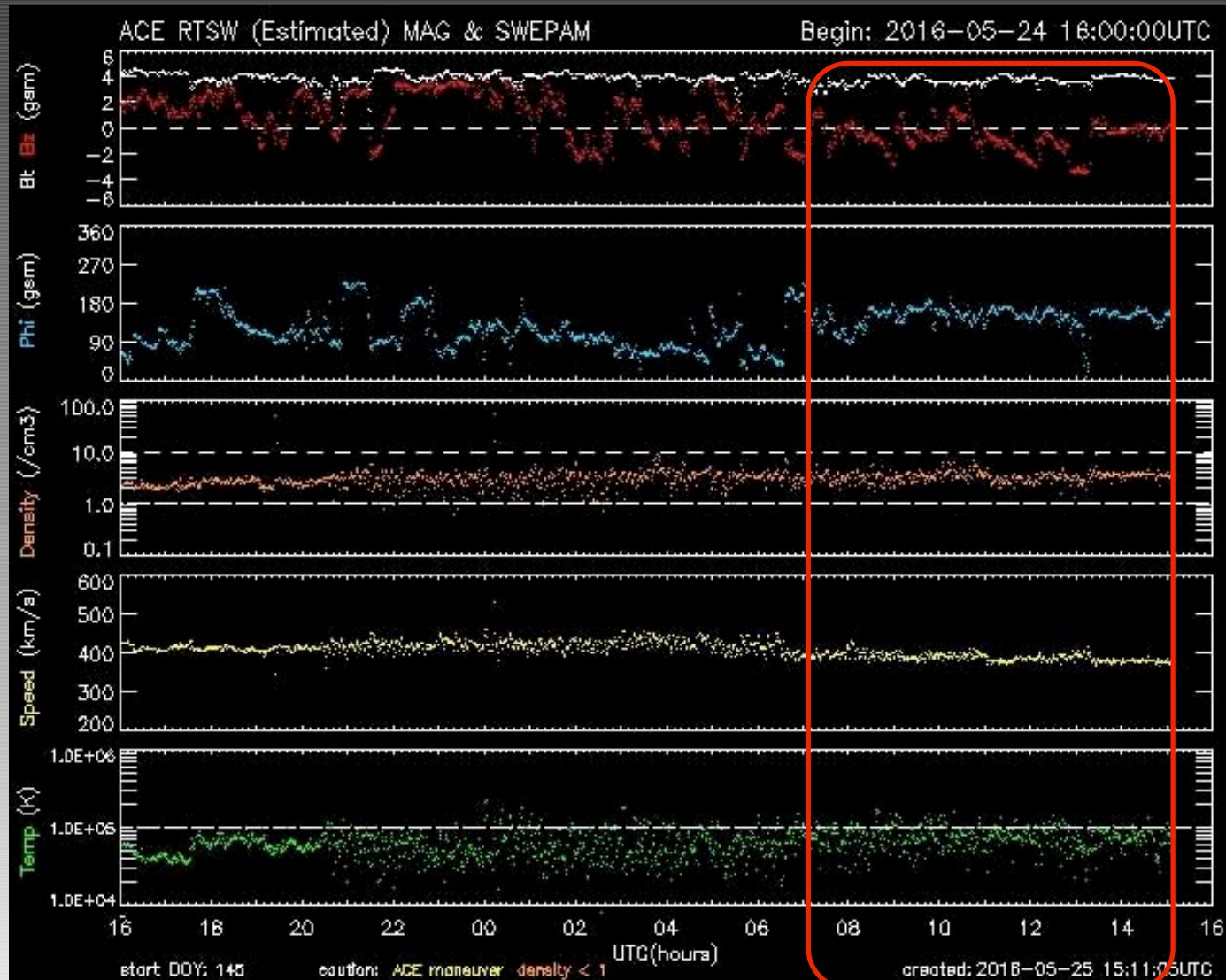


Near an Earth orbit, the **spiral angle** is

$$\tan \theta = \frac{-v_{\phi}^{SW}}{v_r^{SW}} \approx \frac{(1 \text{ AU}) \Omega_{\odot}}{v_r^{SW}} \sim \frac{1.5 \times 10^8 \text{ km} \times (2.87 \times 10^{-6} \text{ rad s}^{-1})}{400 \text{ km s}^{-1}} = 1.07$$

$$\Rightarrow \theta \sim 47^{\circ} \quad (\phi^{GSM} \sim 133^{\circ})$$

# Solar wind observed near an Earth orbit



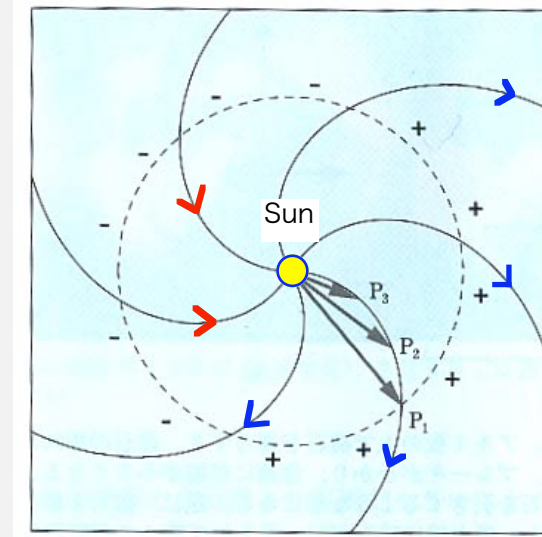
nearly steady state

# Sector structure of IMF

There are **inward IMF area** and **outward IMF area**.

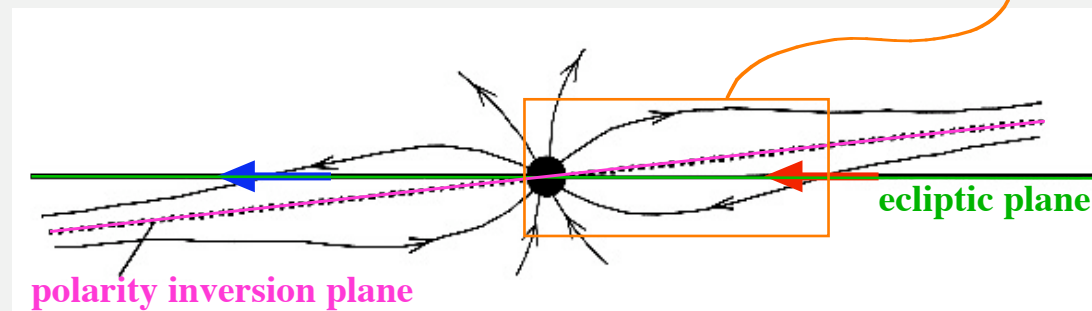


**Sector structure**



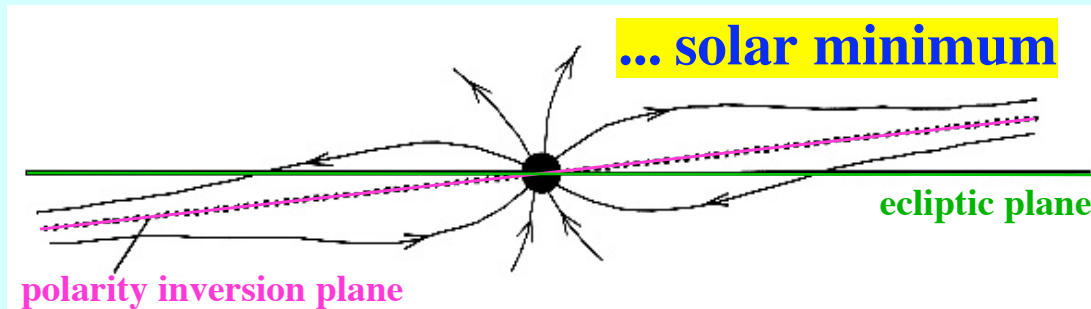
## Origin of sector structure:

**Polarity inversion plane is not coincident with ecliptic plane.**

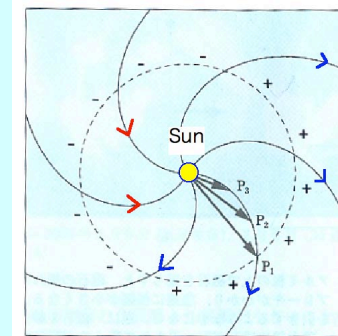


# Time variation of the sector structure...

## Two-sector structure

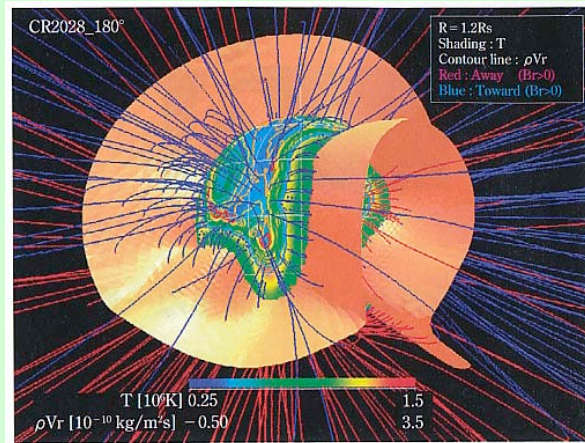


... solar minimum



Polarity inversion plane is almost flat.

## Multi-sector structure



... solar maximum

Polarity inversion plane is strongly rippled because a number of solar active regions exist.

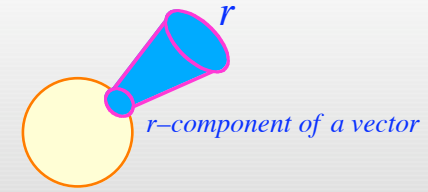
From T. Tanaka

# Solar wind model

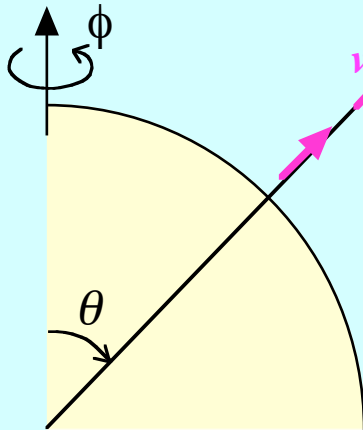
# Hydrodynamic model

# Parker's 1-dimensional model

(spherically symmetric, isothermal, steady, no rotation)  
(depends on  $r$ ;  $r$ -component of a vector is considered)



$(r, \theta, \phi)$ ... spherical coordinates



steady  $\Rightarrow \frac{\partial}{\partial t} = 0$  (time-independent)

1-dimensional model  $\Rightarrow$  all quantities depend only on  $r$

spherically symmetric (isotropic) outflow

$\Rightarrow$  flow velocity has only a radial component ( $v_r$ )

$$\mathbf{v}(r, \theta, \phi) = v_r(r, \theta, \phi) \hat{\mathbf{r}} + v_\theta(r, \theta, \phi) \hat{\boldsymbol{\theta}} + v_\phi(r, \theta, \phi) \hat{\boldsymbol{\phi}}$$

$\longrightarrow v_r(r) \hat{\mathbf{r}}$

isothermal

$\Rightarrow T = \text{const.}$

thermal conduction efficiently works to make temperature uniform

## Basic equations:

**Mass conservation..**  $\rho v_r A = \text{const.}$

spherically symmetric  $\Rightarrow A(r) \propto r^2$

**Momentum equation...**

$$\rho v_r \frac{\partial v_r}{\partial r} = - \frac{\partial p}{\partial r} - \frac{G M_\odot \rho}{r^2}$$

**Energy equation...**

$$T = \text{const.} \left( \frac{p}{\rho} = \frac{k_B T}{\bar{m}} = v_c^2 = \text{const.} \quad v_c = \sqrt{\frac{k_B T}{\bar{m}}} \right)$$

isothermal sound speed

eliminate  $\rho$  and  $p$

$$\left( v_r - \frac{v_c^2}{v_r} \right) \frac{\partial v_r}{\partial r} = \frac{2 v_c^2}{r} - \frac{G M_\odot}{r^2}$$

