

Physical properties of shear Alfvén wave (SAW): (obtained from amplitude relations)

Four vector relation (k , B_0 , v^* , B^*)

- $\bullet k \cdot B^* = 0 (\nabla \cdot B_1 = 0) \Rightarrow k \perp B^*, k \cdot v^* = 0 (\nabla \cdot v_1 = 0) \Rightarrow k \perp v^* \dots \text{transverse wave (incompressive)}$
 - $\bullet \omega B^* = -k \times [v^* \times B_0] = -[(k \cdot B_0) v^* - \underline{(k \cdot v^*) B_0}] = -(k \cdot B_0) v^*$
 $\qquad\qquad\qquad = 0$

$$\mathbf{v}^* = -\frac{\omega}{(\mathbf{k} \cdot \mathbf{B}_0)} \mathbf{B}^* \Rightarrow \begin{cases} \text{phase difference of } \pi \text{ for } \mathbf{k} \cdot \mathbf{B}_0 > 0 \text{ (forward propagation)} \\ \text{same phase for } \mathbf{k} \cdot \mathbf{B}_0 < 0 \text{ (backward propagation)} \end{cases}$$

Oscillation directions of velocity and magnetic field are aligned: $v^* = \begin{cases} - & \text{(forward propagation)} \\ + & \text{(backward propagation)} \end{cases} \frac{\mathbf{B}^*}{\sqrt{\mu_0 \rho_{\text{gas}}}}$

They are perpendicular to B_0

- $v^* \bullet B_0 = 0$, $v^* = \pm \frac{B^*}{\sqrt{\mu_0 \rho_0}} \Rightarrow B^* \bullet B_0 = 0$

... common property

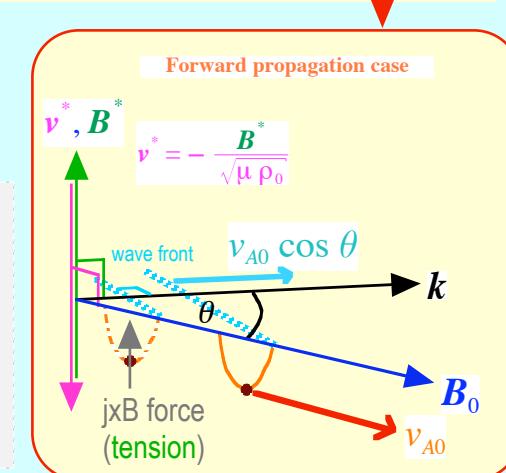
- Driving force:

$$\mathbf{j}^* \times \mathbf{B}_0 = \frac{1}{\mu_0} [\mathbf{k} \times \mathbf{B}^*] \times \mathbf{B}_0 = \boxed{(\mathbf{B}_0 \cdot \mathbf{k}) \frac{\mathbf{B}^*}{\mu_0}} - \boxed{\left(\frac{\mathbf{B}^* \cdot \mathbf{B}_0}{\mu_0} \right) \mathbf{k}}$$

0 except $\theta = \pi/2$ **= 0**

Magnetic tension is the driving force of SAW.

$$p_m^* = \frac{2 \mathbf{B}^* \cdot \mathbf{B}_0}{2 \mu_0}$$



2) $\mathbf{k} \cdot \mathbf{v}^* \neq 0 \Rightarrow \omega^2 - k^2 v_{A0}^2 = 0 \dots$ similar to sound wave (compressive)

=> Compressional Alfvén wave

Velocity diagram

Dispersion relation of compressional Alfvén wave: $\omega(k) = \pm k v_{A0}$

Phase velocity

$$v_p = \frac{\omega}{k} \hat{k} = \pm v_{A0} \hat{k} \quad (+\text{... forward}, -\text{... backward})$$

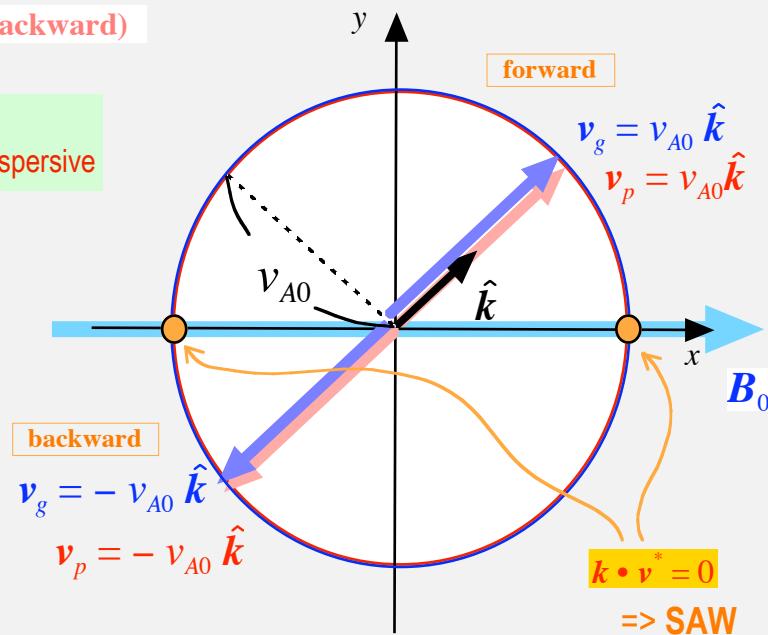
$|v_p|$ is independent of the direction of \mathbf{k} (\hat{k})... isotropic

$|v_p|$ is independent of the magnitude of \mathbf{k} ($|k|$)... non-dispersive

Group velocity

$$v_g = \frac{\partial \omega}{\partial k} = \pm v_{A0} \hat{k}$$

Same speed in all directions
Equal to the phase velocity



Physical properties of compressional Alfvén wave (CAW):

(obtained from amplitude relations)

Four vector relation ($k, \mathbf{B}_0, \mathbf{v}^*, \mathbf{B}^*$)

$$\bullet k \cdot \mathbf{B}^* = 0 \Rightarrow \mathbf{k} \perp \mathbf{B}^*, \quad \mathbf{v}^* \cdot \mathbf{B}_0 = 0 \Rightarrow \mathbf{v}^* \perp \mathbf{B}_0 \quad \text{... common property}$$

(div $\mathbf{B} = 0$)

$$\bullet \omega \rho_0 \mathbf{v}^* = -\frac{1}{\mu_0} [\mathbf{k} \times \mathbf{B}^*] \times \mathbf{B}_0 = -\frac{1}{\mu_0} [(k \cdot \mathbf{B}_0) \mathbf{B}^* - (\mathbf{B}^* \cdot \mathbf{B}_0) \mathbf{k}]$$

$\Rightarrow \mathbf{v}^*$ is on the plane formed by \mathbf{B}^* and \mathbf{k} ($\mathbf{k} \perp \mathbf{B}^*$)

$$\bullet \omega \mathbf{B}^* = -\mathbf{k} \times [\mathbf{v}^* \times \mathbf{B}_0] = -[(k \cdot \mathbf{B}_0) \mathbf{v}^* - (k \cdot \mathbf{v}^*) \mathbf{B}_0]$$

$\neq 0$

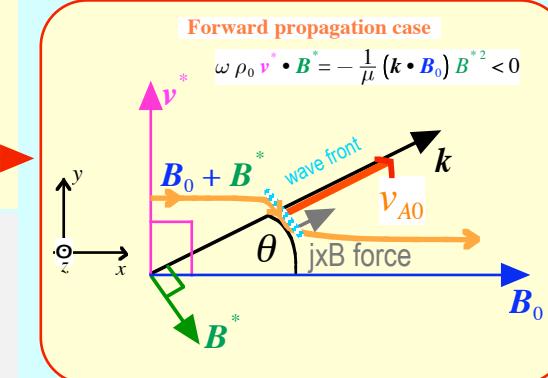
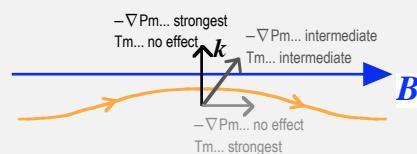
$\Rightarrow \mathbf{B}^*$ is on the plane formed by \mathbf{v}^* and \mathbf{B}_0 ($\mathbf{v}^* \perp \mathbf{B}_0$)

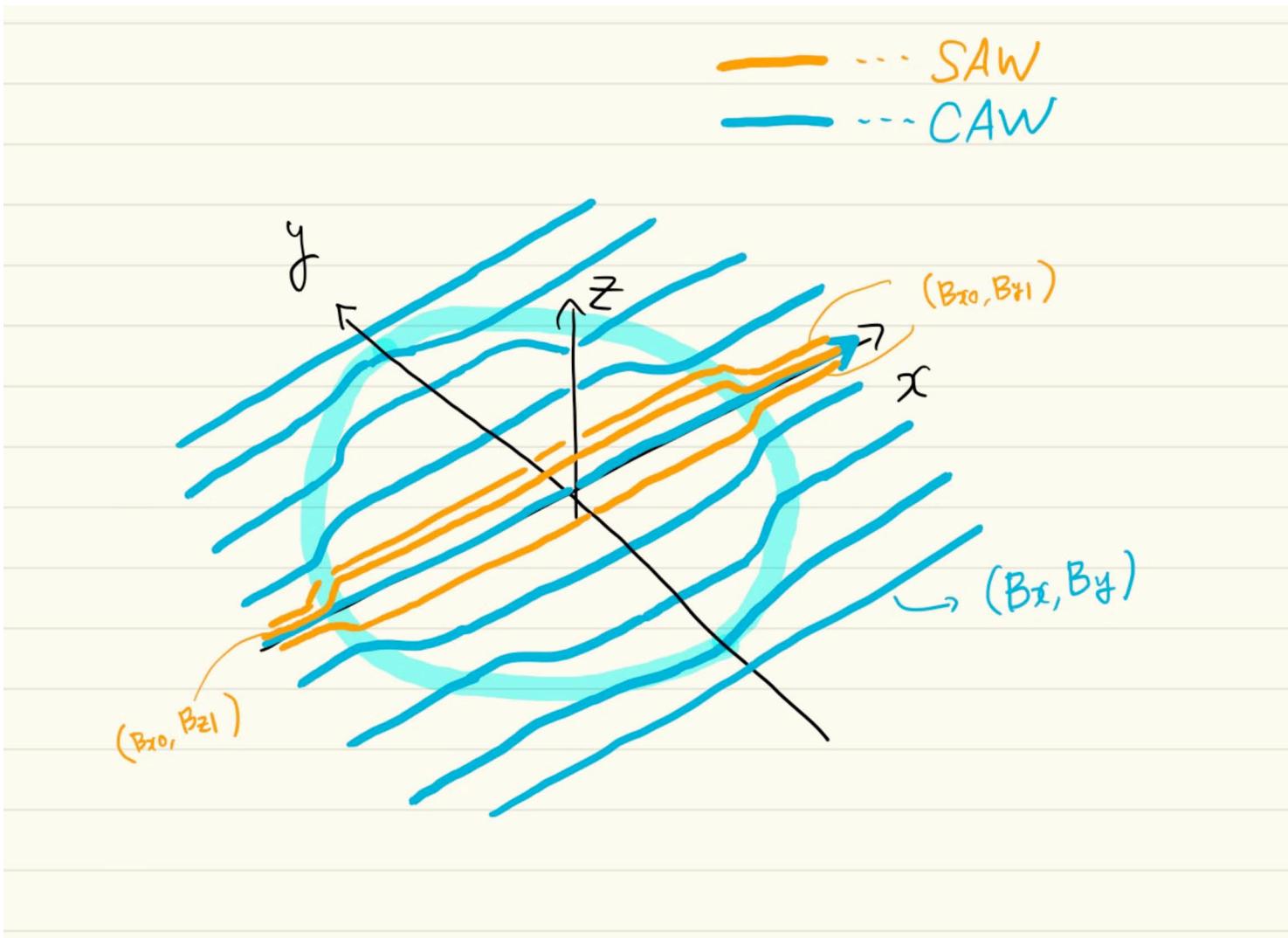
Four vectors $\mathbf{v}^*, \mathbf{B}^*, \mathbf{k}, \mathbf{B}_0$ are on the same plane.

Driving force:

$$\mathbf{j}^* \times \mathbf{B}_0 = \frac{1}{\mu_0} [\mathbf{k} \times \mathbf{B}^*] \times \mathbf{B}_0 = \underbrace{(\mathbf{B}_0 \cdot \mathbf{k}) \frac{\mathbf{B}^*}{\mu_0}}_{\neq 0 \text{ except } \theta = \pm\pi/2} - \underbrace{\left(\frac{\mathbf{B}^* \cdot \mathbf{B}_0}{\mu_0} \right) \mathbf{k}}_{\neq 0 \text{ except } \theta = 0, \pi}$$

Both **tension force** and **magnetic pressure gradient force** are the driving force of CAW.





Magnetoacoustic wave

$p_0 \neq 0, p_1 \neq 0, B_0 \neq 0, B_1 \neq 0$

Linearized ideal MHD equations

$$\begin{cases} \frac{\partial \rho_1(x, y, z, t)}{\partial t} = -\rho_0 \nabla \cdot \mathbf{v}_1(x, y, z, t) \\ \rho_0 \frac{\partial \mathbf{v}_1(x, y, z, t)}{\partial t} = -\nabla p_1(x, y, z, t) + \frac{1}{\mu_0} (\nabla \times \mathbf{B}_1(x, y, z, t)) \times \mathbf{B}_0 \\ \frac{\partial p_1(x, y, z, t)}{\partial t} = \gamma \frac{p_0}{\rho_0} \frac{\partial \rho_1(x, y, z, t)}{\partial t} \\ \frac{\partial \mathbf{B}_1(x, y, z, t)}{\partial t} = \nabla \times (\mathbf{v}_1(x, y, z, t) \times \mathbf{B}_0) \end{cases}$$

Partial differential equations

$$\begin{aligned} \rho_1(x, y, z, t) &= \rho^* e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} \\ \mathbf{v}_1(x, y, z, t) &= \mathbf{v}^* e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} \\ p_1(x, y, z, t) &= p^* e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} \\ \mathbf{B}_1(x, y, z, t) &= \mathbf{B}^* e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} \end{aligned}$$

$$\begin{aligned} -i\omega\rho^*e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} &= -\rho_0 i \mathbf{k} \cdot \mathbf{v}^* e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} \\ -i\omega\rho_0\mathbf{v}^*e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} &= -i\mathbf{k} p^* e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} + \frac{1}{\mu_0} [i\mathbf{k} \times \mathbf{B}^* e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)}] \times \mathbf{B}_0 \\ -i\omega p^*e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} &= -\gamma \frac{p_0}{\rho_0} i\omega\rho^*e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} \\ -i\omega\mathbf{B}^*e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} &= i\mathbf{k} \times [\mathbf{v}^* e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} \times \mathbf{B}_0] \end{aligned}$$

Algebraic equations

$$\begin{aligned} \omega\rho^* &= \rho_0 \mathbf{k} \cdot \mathbf{v}^* \\ \omega\rho_0\mathbf{v}^* &= \mathbf{k} p^* - \frac{1}{\mu_0} (\mathbf{k} \times \mathbf{B}^*) \times \mathbf{B}_0 \\ \omega p^* &= \gamma \frac{p_0}{\rho_0} \omega\rho^* \\ \omega\mathbf{B}^* &= -\mathbf{k} \times (\mathbf{v}^* \times \mathbf{B}_0) \end{aligned}$$

$$\omega \rho^* = \rho_0 \mathbf{k} \cdot \mathbf{v}^*$$

$$\omega \rho_0 \mathbf{v}^* = \mathbf{k} p^* - \frac{1}{\mu_0} (\mathbf{k} \times \mathbf{B}^*) \times \mathbf{B}_0 = \mathbf{k} p^* - \frac{1}{\mu_0} [(\mathbf{k} \cdot \mathbf{B}_0) \mathbf{B}^* - (\mathbf{B}^* \cdot \mathbf{B}_0) \mathbf{k}]$$

$$\omega p^* = \gamma \frac{p_0}{\rho_0} \omega \rho^*$$

$$\omega \mathbf{B}^* = -\mathbf{k} \times (\mathbf{v}^* \times \mathbf{B}_0) = -[(\mathbf{k} \cdot \mathbf{B}_0) \mathbf{v}^* - (\mathbf{k} \cdot \mathbf{v}^*) \mathbf{B}_0]$$

$$\nabla \cdot \mathbf{B}_1 = 0$$

oscillation direction of magnetic field \perp propagation direction

$$\bullet \mathbf{B}_0$$

$$(\omega \rho_0 \mathbf{v}^*) \bullet \mathbf{B}_0 = (\mathbf{k} \cdot \mathbf{B}_0) p^* - \left(\frac{1}{\mu_0} [\mathbf{k} \times \mathbf{B}^*] \times \mathbf{B}_0 \right) \bullet \mathbf{B}_0 \xrightarrow{= 0} \omega \rho_0 (\mathbf{v}^* \bullet \mathbf{B}_0) = (\mathbf{k} \cdot \mathbf{B}_0) p^*$$

$$\bullet \mathbf{k}$$

$$(\omega \rho_0 \mathbf{v}^*) \bullet \mathbf{k} = (\mathbf{k} \cdot \mathbf{k}) p^* - \left(\frac{1}{\mu_0} [(\mathbf{k} \cdot \mathbf{B}_0) \mathbf{B}^* - (\mathbf{B}^* \cdot \mathbf{B}_0) \mathbf{k}] \right) \bullet \mathbf{k}$$

$$\Rightarrow \omega \rho_0 (\mathbf{k} \cdot \mathbf{v}^*) = k^2 p^* - \frac{1}{\mu_0} [(\mathbf{k} \cdot \mathbf{B}_0) (\mathbf{k} \cdot \mathbf{B}^*) - (\mathbf{B}^* \cdot \mathbf{B}_0) k^2] \xrightarrow{= 0}$$

$$\bullet \mathbf{B}_0$$

$$\omega \mathbf{B}^* \bullet \mathbf{B}_0 = -[(\mathbf{k} \cdot \mathbf{B}_0) \mathbf{v}^* - (\mathbf{k} \cdot \mathbf{v}^*) \mathbf{B}_0] \bullet \mathbf{B}_0$$

$$\Rightarrow \omega \mathbf{B}^* \bullet \mathbf{B}_0 = -[(\mathbf{k} \cdot \mathbf{B}_0) \mathbf{v}^* \bullet \mathbf{B}_0 - (\mathbf{k} \cdot \mathbf{v}^*) B_0^2] \xrightarrow{\quad} \omega \mathbf{B}^* \bullet \mathbf{B}_0 = -[(\mathbf{k} \cdot \mathbf{B}_0)^2 \frac{p^*}{\omega \rho_0} - (\mathbf{k} \cdot \mathbf{v}^*) B_0^2]$$

$$\bullet \mathbf{k}$$

$$0 = 0$$

$$\mathbf{v}^* \bullet \mathbf{B}_0 = \frac{\mathbf{k} \cdot \mathbf{B}_0}{\omega \rho_0} p^*$$

$$\omega \rho^* = \rho_0 k \cdot v^* \Rightarrow \rho^* = \frac{\rho_0}{\omega} k \cdot v^*$$

$$\omega p^* = \gamma \frac{p_0}{\rho_0} \omega \rho^* \Rightarrow p^* = \gamma \frac{p_0}{\rho_0} \rho^*$$

$$\omega \rho_0 (k \cdot v^*) = k^2 p^* + \frac{1}{\mu_0} (B^* \cdot B_0) k^2$$

$$\omega B^* \cdot B_0 = - \left[(k \cdot B_0)^2 \frac{p^*}{\omega \rho_0} - (k \cdot v^*) B_0^2 \right]$$

gas pressure vs. magnetic pressure

$$\left(\frac{\omega^2}{\gamma \frac{p_0}{\rho_0}} - k^2 \right) p^* = \frac{1}{\mu_0} (B^* \cdot B_0) k^2$$

$$p_m^* = \frac{2 B^* \cdot B_0}{2 \mu_0}$$

$$p^* = \gamma \frac{p_0}{\omega} k \cdot v^*$$

eliminate $B^* \cdot B_0$

eliminate p^*

$$\left(\omega^2 - \frac{B_0^2}{\mu_0 \rho_0} k^2 \right) (k \cdot v^*) = \left[\omega k^2 - \frac{k^2}{\omega} \frac{(k \cdot B_0)^2}{\mu_0 \rho_0} \right] \frac{p^*}{\rho_0}$$

$$\left[\omega^4 - \left(\gamma \frac{p_0}{\rho_0} + \frac{B_0^2}{\mu_0 \rho_0} \right) k^2 \omega^2 + \gamma \frac{p_0}{\rho_0} \frac{(k \cdot B_0)^2}{\mu_0 \rho_0} k^2 \right] (k \cdot v^*) = 0$$

$$\left[\omega^4 - \left(\gamma \frac{p_0}{\rho_0} + \frac{B_0^2}{\mu_0 \rho_0} \right) k^2 \omega^2 + \gamma \frac{p_0}{\rho_0} \frac{(k \cdot \mathbf{B}_0)^2}{\mu_0 \rho_0} k^2 \right] (k \cdot v^*) = 0$$

- 1) $k \cdot v^* = 0$... **incompressive**
- 2) $k \cdot v^* \neq 0$... **compressive**

1) $\mathbf{k} \cdot \mathbf{v}^* = 0 \Rightarrow \nabla \cdot \mathbf{v}_1 = 0 \dots$ incompressional wave \Rightarrow Shear Alfvén wave

$\omega \rho^* = \rho_0 \mathbf{k} \cdot \mathbf{v}^*$

$\omega \rho_0 \mathbf{v}^* = \mathbf{k} p^* - \frac{1}{\mu_0} (\mathbf{k} \times \mathbf{B}^*) \times \mathbf{B}_0 = \mathbf{k} p^* - \frac{1}{\mu_0} [(\mathbf{k} \cdot \mathbf{B}_0) \mathbf{B}^* - (\mathbf{B}^* \cdot \mathbf{B}_0) \mathbf{k}]$

$\omega p^* = \gamma \frac{p_0}{\rho_0} \omega \rho^*$

$\omega \mathbf{B}^* = -\mathbf{k} \times (\mathbf{v}^* \times \mathbf{B}_0) = -[(\mathbf{k} \cdot \mathbf{B}_0) \mathbf{v}^* - (\mathbf{k} \cdot \mathbf{v}^*) \mathbf{B}_0] = 0$

$\mathbf{v}^* \cdot \mathbf{B}_0 = \frac{\mathbf{k} \cdot \mathbf{B}_0}{\omega \rho_0} p^* = 0$

$\rho^* = 0$

$p^* = \gamma \frac{p_0}{\rho_0} \quad \rho^* = 0$

$\omega \rho_0 \mathbf{v}^* = -\frac{1}{\mu_0} [(\mathbf{k} \cdot \mathbf{B}_0) \mathbf{B}^* - (\mathbf{B}^* \cdot \mathbf{B}_0) \mathbf{k}]$

$\omega \mathbf{B}^* = -(\mathbf{k} \cdot \mathbf{B}_0) \mathbf{v}^*$

eliminate \mathbf{B}^*

$\omega^2 \rho_0 \mathbf{v}^* = -\frac{1}{\mu_0} \left[-(\mathbf{k} \cdot \mathbf{B}_0)^2 \mathbf{v}^* + (\mathbf{k} \cdot \mathbf{v}^*) (\mathbf{k} \cdot \mathbf{B}_0) \mathbf{B}_0 + [(\mathbf{k} \cdot \mathbf{B}_0) (\mathbf{v}^* \cdot \mathbf{B}_0) - (\mathbf{k} \cdot \mathbf{v}^*) B_0^2] \mathbf{k} \right] = 0$

$\left[\omega^2 - \frac{(\mathbf{k} \cdot \mathbf{B}_0)^2}{\mu_0 \rho_0} \right] \mathbf{v}^* = \mathbf{0} \Rightarrow \omega = \pm \frac{\mathbf{k} \cdot \mathbf{B}_0}{\sqrt{\mu_0 \rho_0}} \dots$ dispersion relation of SAW

Velocity diagram

Dispersion relation of SAW: $\omega(k) = \pm \frac{\mathbf{k} \cdot \mathbf{B}_0}{\sqrt{\mu_0 \rho_0}} = \pm k v_{A0} \cos \theta$

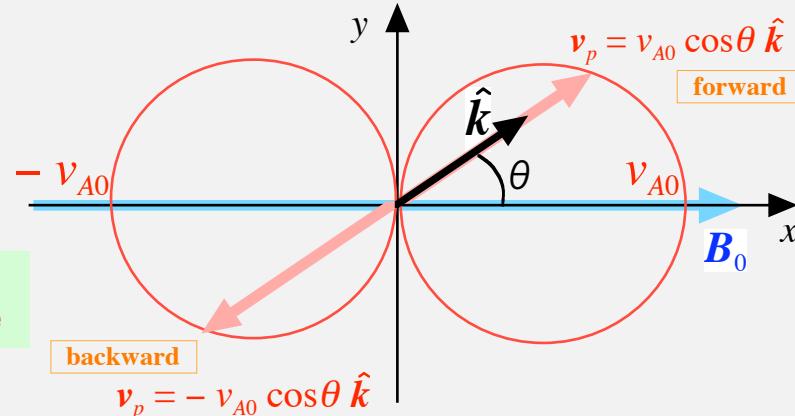
Phase velocity

$$\mathbf{v}_p = \frac{\omega}{k} \hat{\mathbf{k}} = \pm v_{A0} \cos \theta \hat{\mathbf{k}}$$

(... forward, ... backward)

$|\mathbf{v}_p|$ is dependent on the direction of \mathbf{k} ($\hat{\mathbf{k}}$)... non-isotropic

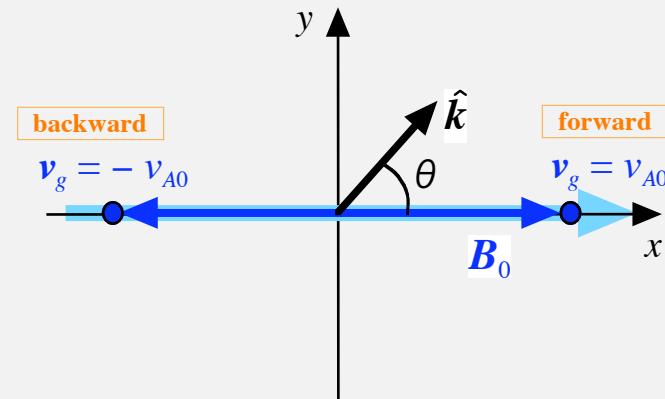
$|\mathbf{v}_p|$ is independent of the magnitude of \mathbf{k} ($|k|$)... non-dispersive



Consider why shear Alfvén wave is not affected by gas pressure.

Group velocity

$$\begin{aligned} \mathbf{v}_g &= \frac{\partial \omega}{\partial \mathbf{k}} = \pm \left(\frac{\partial}{\partial k_x} \frac{\partial}{\partial k_y} \frac{\partial}{\partial k_z} \right) \frac{\mathbf{k} \cdot \mathbf{B}_0}{\sqrt{\mu_0 \rho_0}} \\ &= \pm \left(\frac{\partial}{\partial k_x} \frac{\partial}{\partial k_y} \frac{\partial}{\partial k_z} \right) \frac{k_x B_0 + k_y 0 + k_z 0}{\sqrt{\mu_0 \rho_0}} = \pm \begin{pmatrix} \frac{B_0}{\sqrt{\mu_0 \rho_0}} \\ 0 \\ 0 \end{pmatrix} = \pm \begin{pmatrix} v_{A0} \\ 0 \\ 0 \end{pmatrix} \end{aligned}$$



Energy (wave packet) is always transported along magnetic field \mathbf{B}_0 .

$$2) \quad \mathbf{k} \cdot \mathbf{v}^* \neq 0 \Rightarrow \omega^4 - \left(\gamma \frac{p_0}{\rho_0} + \frac{B_0^2}{\mu_0 \rho_0} \right) k^2 \omega^2 + \gamma \frac{p_0}{\rho_0} \frac{(\mathbf{k} \cdot \mathbf{B}_0)^2}{\mu_0 \rho_0} k^2 = 0$$

compressional wave



$$\omega^4 - (c_{s0}^2 + v_{A0}^2) k^2 \omega^2 + c_{s0}^2 v_{A0}^2 k^4 \cos^2 \theta = 0$$



When θ is zero (propagation along magnetic field), this is reduced to

$$(\omega^2 - k^2 c_{s0}^2)(\omega^2 - k^2 v_{A0}^2) = 0,$$

so sound wave and shear Alfvén wave are completely decoupled.

Consider this from a viewpoint of driving force.

Two types of compressional MHD wave: **slow wave & fast wave**

$$\omega(\mathbf{k}) = \pm k \sqrt{\frac{c_{s0}^2 + v_{A0}^2}{2}} \pm \frac{\sqrt{c_{s0}^4 + v_{A0}^4 - 2 c_{s0}^2 v_{A0}^2 \cos 2\theta}}{2}$$

forward (+)
 backward (-)

fast wave (+)
 slow wave (-)

$\cos 2\theta = 2 \cos^2 \theta - 1$

Slow wave

Dispersion relation: $\omega^2 = k^2 \left(\frac{c_{s0}^2 + v_{A0}^2}{2} - \frac{\sqrt{c_{s0}^4 + v_{A0}^4 - 2 c_{s0}^2 v_{A0}^2 \cos 2\theta}}{2} \right) \leq k^2 c_{s0}^2$ Confirm this.

Gas pressure and magnetic pressure work in an uncooperative way.

gas pressure vs. magnetic pressure

$$\left(\frac{\omega^2}{\gamma \frac{p_0}{\rho_0}} - k^2 \right) p^* = \frac{1}{\mu_0} (B^* \cdot B_0) k^2$$

1 out of phase

$$\frac{\omega^2}{c_{s0}^2} - k^2 \leq 0$$

Phase velocity

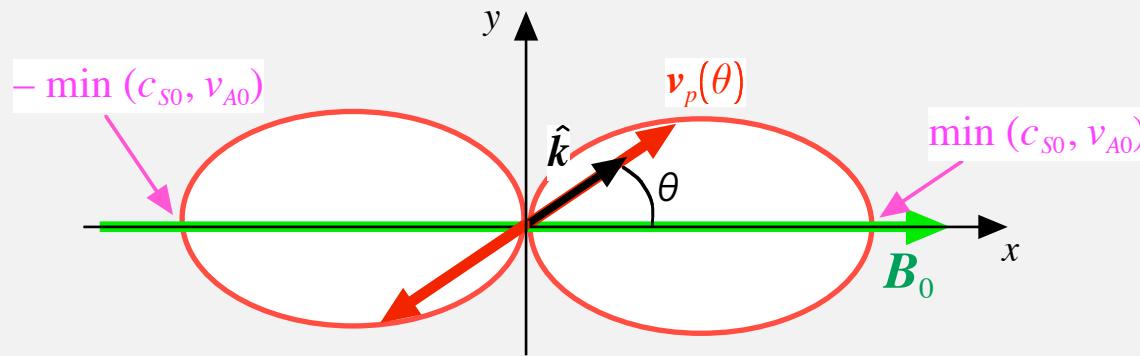
$$v_p = \frac{\omega(k, \theta)}{k} \hat{k} = \pm \sqrt{\frac{c_{s0}^2 + v_{A0}^2}{2} - \frac{\sqrt{c_{s0}^4 + v_{A0}^4 - 2 c_{s0}^2 v_{A0}^2 \cos 2\theta}}{2}} \hat{k}$$

$$= v_p(\theta)$$

only depends on the direction of k

$|v_p|$ is dependent on the direction of k (\hat{k})... non-isotropic

$|v_p|$ is independent of the magnitude of k ($\|k\|$)... non-dispersive



Group velocity

$$v_g = \frac{\partial \omega}{\partial k}, \quad \omega(k, \theta) = \pm k \sqrt{\frac{c_{s0}^2 + v_{A0}^2}{2} - \frac{\sqrt{c_{s0}^4 + v_{A0}^4 - 2 c_{s0}^2 v_{A0}^2 \cos 2\theta}}{2}}$$

$$k = \sqrt{k_x^2 + k_y^2}$$

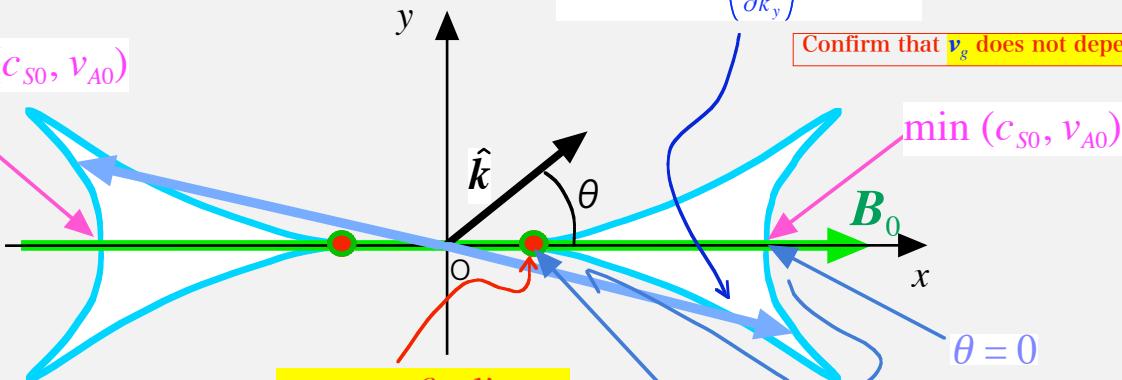
$$\cos(2\theta) = \frac{1 - \tan^2 \theta}{1 + \tan^2 \theta} = \frac{1 - \frac{k_y^2}{k_x^2}}{1 + \frac{k_y^2}{k_x^2}} = \frac{k_x^2 - k_y^2}{k_x^2 + k_y^2}$$

$$\omega(k_x, k_y)$$

$$\mathbf{v}_g(k, \theta) = \begin{pmatrix} \frac{\partial \omega}{\partial k_x} \\ \frac{\partial \omega}{\partial k_y} \end{pmatrix} = \mathbf{v}_g(\theta)$$

$-\min(c_{s0}, v_{A0})$

Confirm that \mathbf{v}_g does not depend on $k = |\mathbf{k}|$.



$$c_T \equiv \frac{c_{s0} v_{A0}}{\sqrt{c_{s0}^2 + v_{A0}^2}}$$

tube velocity

Fast wave

Dispersion relation: $\omega^2 = k^2 \left(\frac{c_{s0}^2 + v_{A0}^2}{2} + \frac{\sqrt{c_{s0}^4 + v_{A0}^4 - 2 c_{s0}^2 v_{A0}^2 \cos 2\theta}}{2} \right) \geq k^2 c_{s0}^2$ Confirm this.

Gas pressure and magnetic pressure work **in a cooperative way**.

gas pressure vs. magnetic pressure

$$\left(\frac{\omega^2}{\gamma P_0} - k^2 \right) p^* = \frac{1}{\mu_0} (\mathbf{B}^* \cdot \mathbf{B}_0) k^2$$

$\frac{\omega^2}{c_{s0}^2} - k^2 \geq 0$

in phase

Phase velocity

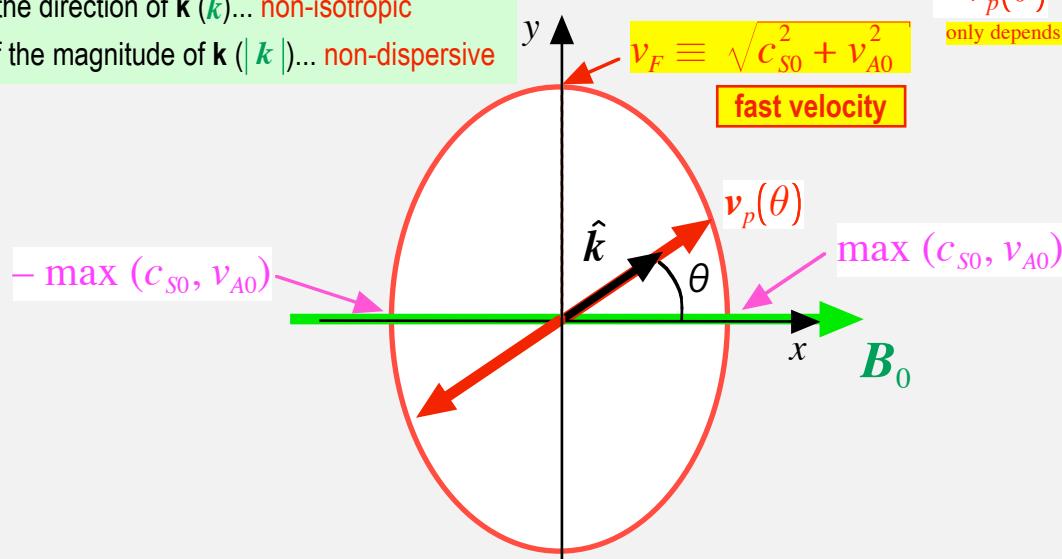
$$\mathbf{v}_p = \frac{\omega(k, \theta)}{k} \hat{\mathbf{k}} = \pm \sqrt{\frac{c_{s0}^2 + v_{A0}^2}{2} + \frac{\sqrt{c_{s0}^4 + v_{A0}^4 - 2 c_{s0}^2 v_{A0}^2 \cos 2\theta}}{2}} \hat{\mathbf{k}}$$

$|\mathbf{v}_p|$ is dependent on the direction of \mathbf{k} ($\hat{\mathbf{k}}$)... **non-isotropic**

$|\mathbf{v}_p|$ is independent of the magnitude of \mathbf{k} ($|k|$)... **non-dispersive**

$$= \mathbf{v}_p(\theta)$$

only depends on the direction of \mathbf{k}



Group velocity

$$v_g = \frac{\partial \omega}{\partial k}, \quad \omega(k, \theta) = \pm k \sqrt{\frac{c_{s0}^2 + v_{A0}^2}{2} + \frac{\sqrt{c_{s0}^4 + v_{A0}^4 - 2 c_{s0}^2 v_{A0}^2 \cos 2\theta}}{2}}$$

$$k = \sqrt{k_x^2 + k_y^2}$$

$$\cos(2\theta) = \frac{1 - \tan^2 \theta}{1 + \tan^2 \theta} = \frac{1 - \frac{k_y^2}{k_x^2}}{1 + \frac{k_y^2}{k_x^2}} = \frac{k_x^2 - k_y^2}{k_x^2 + k_y^2}$$

$$\omega(k_x, k_y)$$

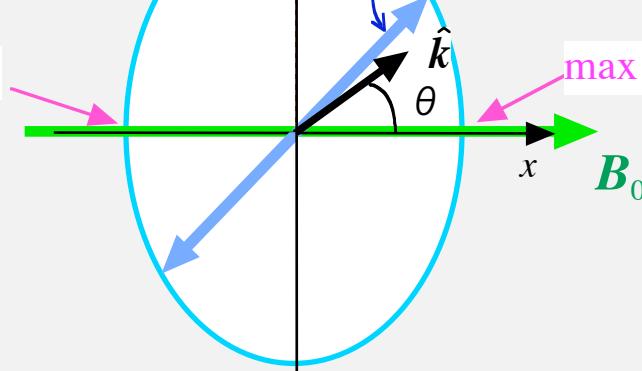
$$v_F \equiv \sqrt{c_{s0}^2 + v_{A0}^2}$$

$$\mathbf{v}_g(k, \theta) = \begin{pmatrix} \frac{\partial \omega}{\partial k_x} \\ \frac{\partial \omega}{\partial k_y} \end{pmatrix} = \mathbf{v}_g(\theta)$$

Confirm that \mathbf{v}_g does not depend on $k = |\mathbf{k}|$.

$$-\max(c_{s0}, v_{A0})$$

$$\max(c_{s0}, v_{A0})$$



Slow wave

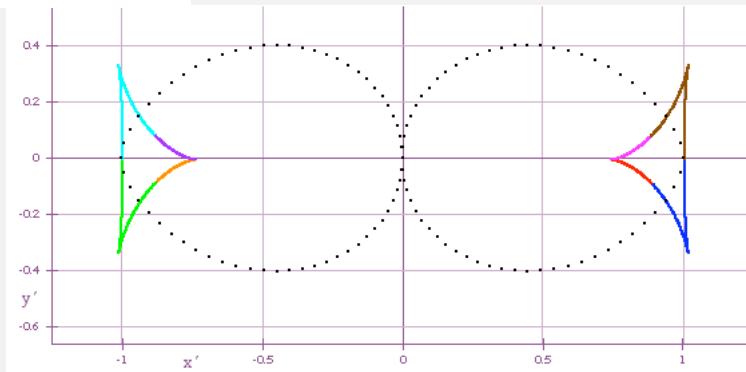
V_p (phase velocity)... dotted

V_g (group velocity)... colored

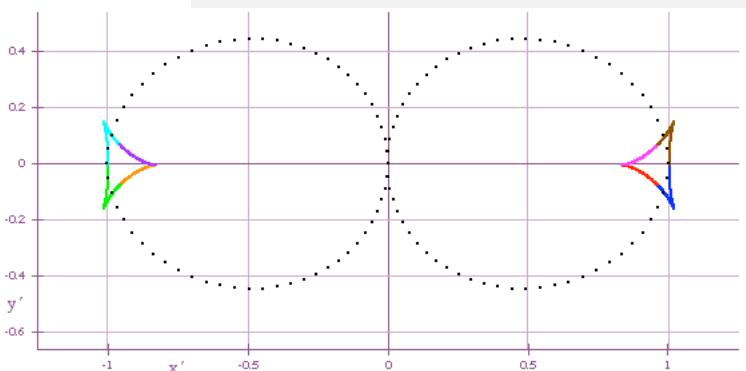
$$0 \leq \theta \leq \frac{\pi}{4} \dots \text{blue}, \frac{\pi}{4} \leq \theta \leq \frac{\pi}{2} \dots \text{red}, \frac{\pi}{2} \leq \theta \leq \frac{3\pi}{4} \dots \text{orange}, \frac{3\pi}{4} \leq \theta \leq \pi \dots \text{green}$$

$$\pi \leq \theta \leq \frac{5\pi}{4} \dots \text{cyan}, \frac{5\pi}{4} \leq \theta \leq \frac{3\pi}{2} \dots \text{purple}, \frac{3\pi}{2} \leq \theta \leq \frac{7\pi}{4} \dots \text{magenta}, \frac{7\pi}{4} \leq \theta \leq 2\pi \dots \text{brown}$$

$$C_S = 1.0, V_A = 1.1$$



$$C_S = 1.0, V_A = 1.5$$



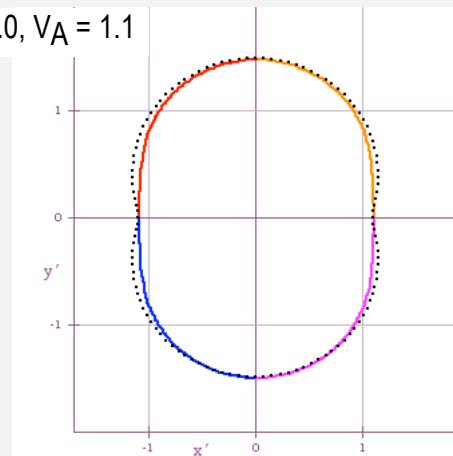
Fast wave

V_p (phase velocity)... dotted

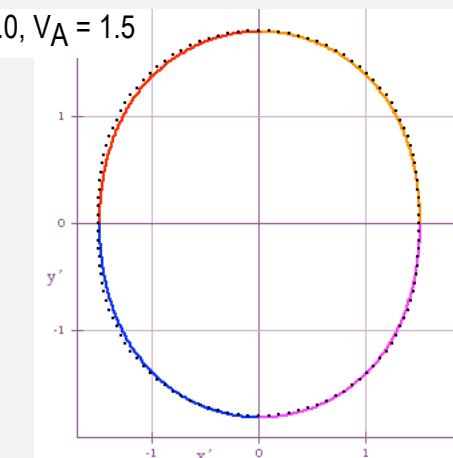
V_g (group velocity)... colored

$$0 \leq \theta \leq \frac{\pi}{2} \dots \text{orange}, \frac{\pi}{2} \leq \theta \leq \pi \dots \text{red}, \pi \leq \theta \leq \frac{3\pi}{2} \dots \text{blue}, \frac{3\pi}{2} \leq \theta \leq 2\pi \dots \text{magenta}$$

$$C_S = 1.0, V_A = 1.1$$

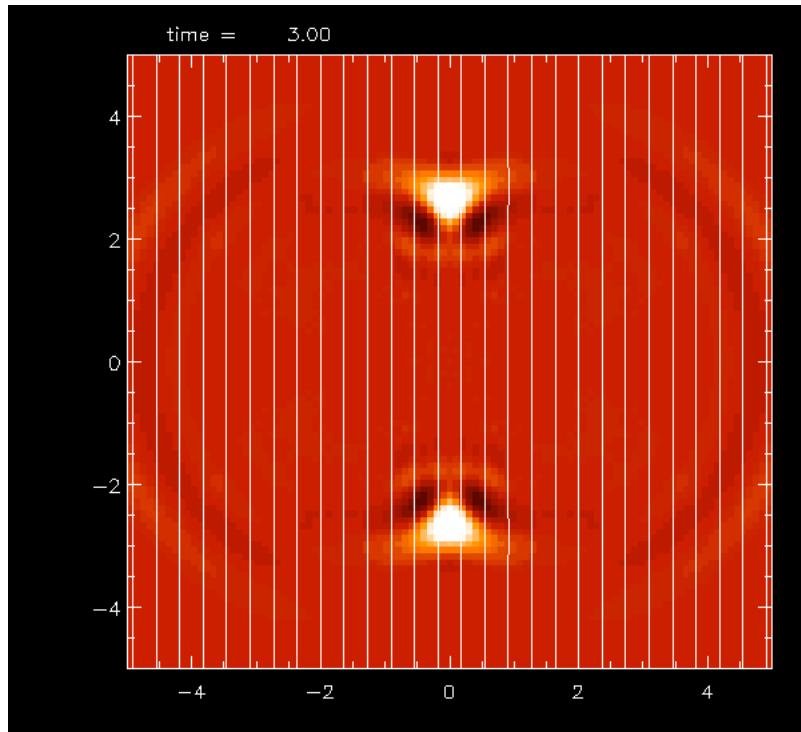


$$C_S = 1.0, V_A = 1.5$$

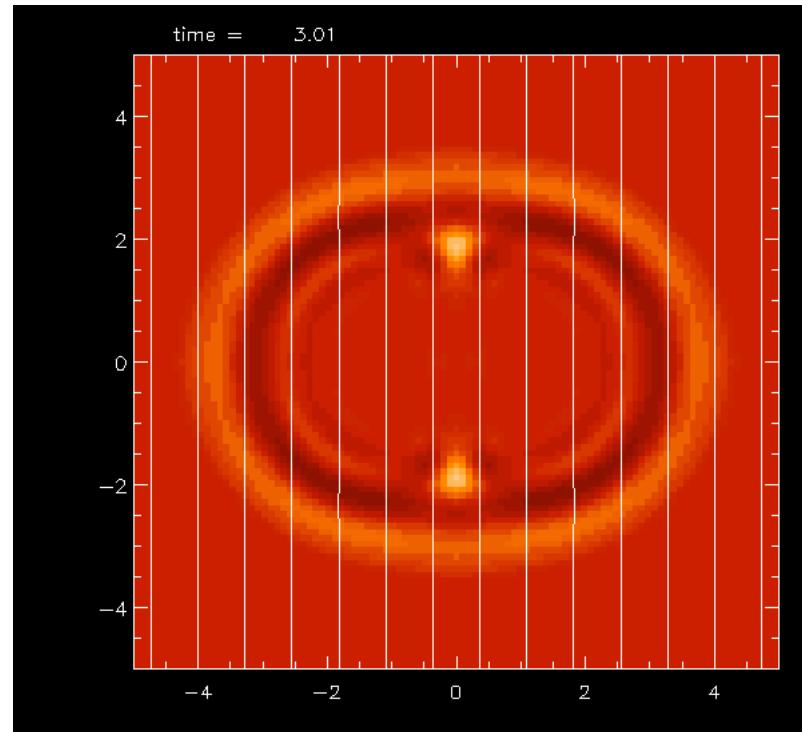


2-dimensional simulation of compressional MHD wave

Low plasma beta case (strong B)



High plasma beta case (weak B)

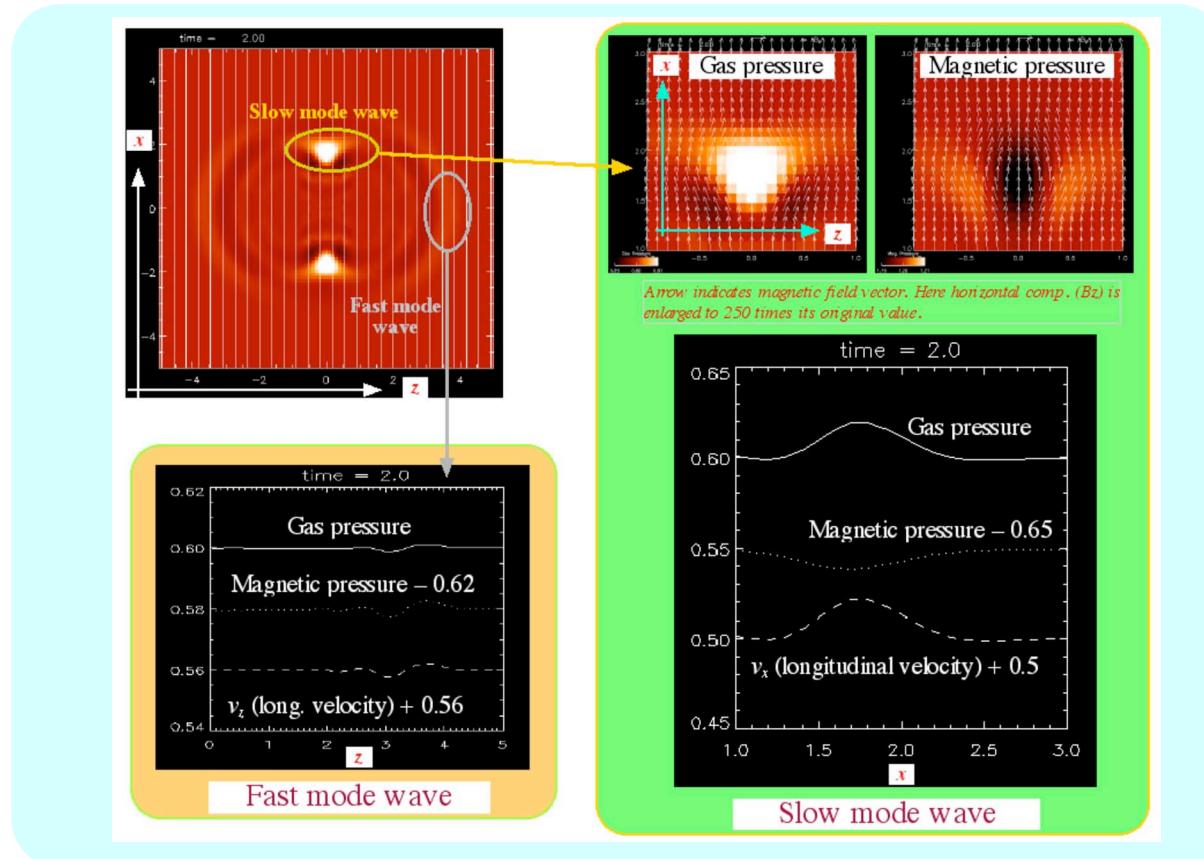


Oscillation continues at the origin for about 1.5-period.

Magnetic field plays a role in **transporting wave energy along it**.

For more details, see the website "http://163.180.179.74/~magara/page31/Topics/MHD_waves/mhdw.html".

Physical properties of compressional MHD wave:



gas pressure vs. magnetic pressure

$$\left(\frac{\omega^2}{c_{s0}^2} - k^2 \right) p^* = \frac{1}{\mu_0} (\mathbf{B}^* \cdot \mathbf{B}_0) k^2$$

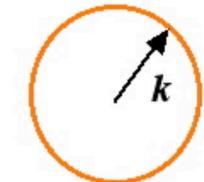
$\frac{\omega}{k} > c_{s0}$ (fast wave) $\Rightarrow p^*$ and p_m^* have positive correlation

p^* and v^* have positive correlation when $\begin{cases} \mathbf{k} \perp \mathbf{B}_0 \text{ for fast-mode wave} \\ \mathbf{k} \parallel \mathbf{B}_0 \text{ for slow-mode wave} \end{cases}$

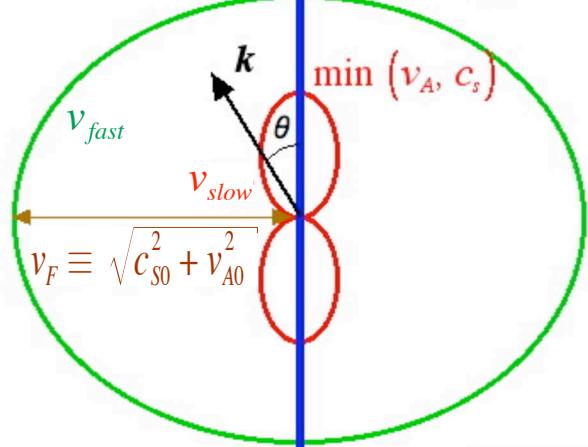
$\frac{\omega}{k} < c_{s0}$ (slow wave) $\Rightarrow p^*$ and p_m^* have negative correlation

Summary of compressional MHD wave

HD wave



$$v_p = v_g = c_s$$

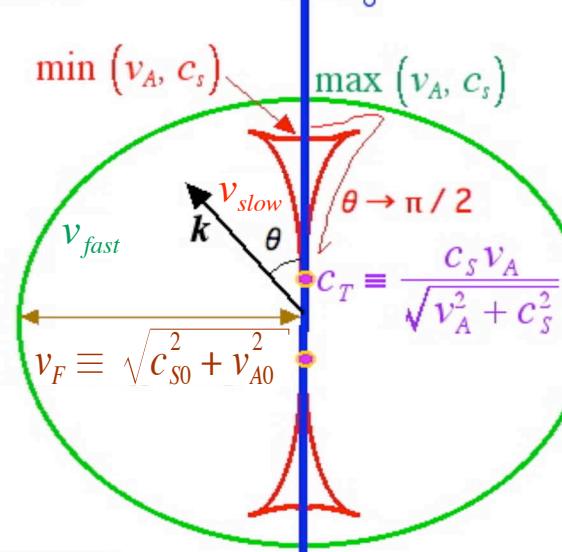


Phase velocity
 $(v_p = \frac{\omega}{k} \hat{k})$

c_s : Sound velocity
 v_A : Alfvén velocity
 c_T : Tube velocity

MHD wave

MHD wave



Group velocity
 $(v_g = \frac{\partial \omega}{\partial k})$

$$c_T < \min(c_{s0}, v_{A0}) < \max(c_{s0}, v_{A0}) < v_F$$