

$$3. \frac{\partial}{\partial t} \left(\frac{p}{\gamma - 1} \right) + \nabla \cdot \left(\frac{p}{\gamma - 1} \mathbf{v} \right) = -p \nabla \cdot \mathbf{v}$$

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$$\frac{p}{\gamma - 1} \nabla \cdot \mathbf{v} + \mathbf{v} \cdot \nabla \left(\frac{p}{\gamma - 1} \right)$$

$$\nabla \cdot \mathbf{v} = -\frac{1}{\rho} \left(\frac{\partial \rho}{\partial t} + \mathbf{v} \cdot \nabla \rho \right)$$

mass conservation

$$\frac{\partial p}{\partial t} + \mathbf{v} \cdot \nabla p - \frac{\gamma p}{\rho} \left[\frac{\partial \rho}{\partial t} + \mathbf{v} \cdot \nabla \rho \right] = 0$$

$$\rho_0 + \rho_1(x, y, z, t)$$

$$\mathbf{0} + \mathbf{v}_1(x, y, z, t)$$

$$p_0 + p_1(x, y, z, t)$$

$$\frac{\partial p}{\partial t} + \mathbf{v} \cdot \nabla p - \frac{\gamma p}{\rho} \left[\frac{\partial \rho}{\partial t} + \mathbf{v} \cdot \nabla \rho \right] = 0$$

$$\frac{\partial \rho_1(x, y, z, t)}{\partial t} + \mathbf{v}_1(x, y, z, t) \cdot \nabla \rho_0 = 0$$

(ρ_0 is constant)

$$\frac{\partial (p_0 + p_1(x, y, z, t))}{\partial t} + \mathbf{v}_1(x, y, z, t) \cdot \nabla (p_0 + p_1(x, y, z, t))$$

2nd-order
 $\square_1 \bigcirc_1 \doteq 0$

$$\frac{\partial p_1(x, y, z, t)}{\partial t}$$

$$\mathbf{v}_1(x, y, z, t) \cdot \nabla p_0 = 0 \quad (p_0 \text{ is constant})$$

$$\gamma \frac{p_0 + p_1(x, y, z, t)}{\rho_0 + \rho_1(x, y, z, t)} = \gamma (p_0 + p_1(x, y, z, t)) (\rho_0 + \rho_1(x, y, z, t))^{-1}$$

$$= \gamma \frac{p_0}{\rho_0} \left(1 + \frac{p_1(x, y, z, t)}{p_0} \right) \left(1 + \frac{\rho_1(x, y, z, t)}{\rho_0} \right)^{-1}$$

$$\approx \gamma \frac{p_0}{\rho_0} \left(1 + \frac{p_1(x, y, z, t)}{p_0} \right) \left(1 - \frac{\rho_1(x, y, z, t)}{\rho_0} + \dots \right)$$

$$\approx \gamma \frac{p_0}{\rho_0} \left(1 + \frac{p_1(x, y, z, t)}{p_0} - \frac{\rho_1(x, y, z, t)}{\rho_0} \right)$$

$(1+x)^{-1} \approx 1-x \text{ when } |x| \ll 1$
 $\left(\frac{p_1(x, y, z, t)}{p_0} \times \frac{\rho_1(x, y, z, t)}{\rho_0} \approx 0 \right)$

$$\frac{\partial p}{\partial t} + \mathbf{v} \cdot \nabla p - \frac{\gamma p}{\rho} \left[\frac{\partial \rho}{\partial t} + \mathbf{v} \cdot \nabla \rho \right] = 0$$



$$\frac{\partial p_1(x, y, z, t)}{\partial t} = \gamma \frac{p_0}{\rho_0} \left(1 + \frac{p_1(x, y, z, t)}{p_0} - \frac{\rho_1(x, y, z, t)}{\rho_0} \right) \frac{\partial \rho_1(x, y, z, t)}{\partial t}$$

Only 1st-order terms remain



$$\frac{\partial p_1(x, y, z, t)}{\partial t} = \gamma \frac{p_0}{\rho_0} \frac{\partial \rho_1(x, y, z, t)}{\partial t}$$

Linearized internal energy equation

$$4. \frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B})$$

$$\begin{matrix} \mathbf{0} + \mathbf{v}_1(x, y, z, t) \\ \mathbf{B}_0 + \mathbf{B}_1(x, y, z, t) \end{matrix} \longrightarrow \frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B})$$

$$\frac{\partial (\mathbf{B}_0 + \mathbf{B}_1(x, y, z, t))}{\partial t} = \nabla \times [\mathbf{v}_1(x, y, z, t) \times (\mathbf{B}_0 + \mathbf{B}_1(x, y, z, t))]$$

$\downarrow$ (green arrow) $\frac{\partial \mathbf{B}_1}{\partial t}$ $\downarrow$ (blue arrow) $\nabla \times (\mathbf{v}_1 \times \mathbf{B}_0)$

$\xrightarrow{\text{2nd-order}} \mathbf{v}_1 \times \mathbf{B}_1 \approx 0$

$\frac{\partial \mathbf{B}_1}{\partial t} = \nabla \times (\mathbf{v}_1 \times \mathbf{B}_0)$ Linearized induction equation

Ideal MHD equations

$$\begin{cases} \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0 \\ \rho \left(\frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} \right) = -\nabla p + \frac{1}{\mu_0} (\nabla \times \mathbf{B}) \times \mathbf{B} \\ \rho \frac{d}{dt} \left(\frac{1}{\gamma - 1} \frac{p}{\rho} \right) + p \nabla \cdot \mathbf{v} = 0 \\ \frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B}) \\ \nabla \cdot \mathbf{B} = 0 \end{cases}$$

$$\begin{aligned} &\rho_0 + \rho_1(x, y, z, t) \\ &\mathbf{0} + \mathbf{v}_1(x, y, z, t) \\ &p_0 + p_1(x, y, z, t) \\ &\mathbf{B}_0 + \mathbf{B}_1(x, y, z, t) \end{aligned}$$

Linearized ideal MHD equations

$$\begin{cases} \frac{\partial \rho_1(x, y, z, t)}{\partial t} = -\rho_0 \nabla \cdot \mathbf{v}_1(x, y, z, t) \\ \rho_0 \frac{\partial \mathbf{v}_1(x, y, z, t)}{\partial t} = -\nabla p_1(x, y, z, t) + \frac{1}{\mu_0} (\nabla \times \mathbf{B}_1(x, y, z, t)) \times \mathbf{B}_0 \\ \frac{\partial p_1(x, y, z, t)}{\partial t} = \gamma \frac{p_0}{\rho_0} \frac{\partial \rho_1(x, y, z, t)}{\partial t} \\ \frac{\partial \mathbf{B}_1(x, y, z, t)}{\partial t} = \nabla \times (\mathbf{v}_1(x, y, z, t) \times \mathbf{B}_0) \\ \nabla \cdot \mathbf{B}_1(x, y, z, t) = 0 \end{cases}$$

Mathematical description of a linear wave (*e^i -form*: $u^* f(\mathbf{k} \cdot \mathbf{r} - \omega t + \phi_0) \Rightarrow |u^*| e^{i\phi_0} e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)}$)
 linear wave solution = combination of $e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)}$ (Fourier decomposition)

$$\rho_1(x, y, z, t) = \rho^* e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)}$$

$$\rho^* = |\rho^*| e^{i\alpha}$$

$$\mathbf{v}_1(x, y, z, t) = \mathbf{v}^* e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)}$$

$$\mathbf{v}_j^* = |\mathbf{v}_j^*| e^{i\beta_j}, j = x, y, z$$

$\alpha, \beta_j, \delta, \varepsilon_j \dots$ initial phase of each quantity

$$p_1(x, y, z, t) = p^* e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)}$$

$$p^* = |p^*| e^{i\delta}$$

$$\text{e.g. } \rho^* e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} = |\rho^*| e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t + \alpha)}$$

$$\mathbf{B}_1(x, y, z, t) = \mathbf{B}^* e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)}$$

$$\mathbf{B}_j^* = |\mathbf{B}_j^*| e^{i\varepsilon_j}, j = x, y, z$$

We always take the **real part** of this form as a linear wave solution:

$$\text{Re} \left[e^{i\phi_0} e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} \right] = \cos(\mathbf{k} \cdot \mathbf{r} - \omega t + \phi_0)$$

As long as **linear term** is considered, we can **always retrieve a correct result by taking the real part**:

$$\text{e.g., } \frac{d}{dx} \cos x = -\sin x$$

$$\frac{d}{dx} e^{ix} = i e^{ix} = i (\cos x + i \sin x) = -\sin x + i \cos x$$

※ *e^i -form* cannot be used for **nonlinear term** (confirm this by considering $(\cos x)^2 \Leftrightarrow (e^{ix})^2$).

Sound wave

$$\mathbf{B}_0 = \mathbf{0}$$

Linearized ideal MHD equations

$$\left\{ \begin{array}{l} \frac{\partial \rho_1(x, y, z, t)}{\partial t} = -\rho_0 \nabla \cdot \mathbf{v}_1(x, y, z, t) \\ \rho_0 \frac{\partial \mathbf{v}_1(x, y, z, t)}{\partial t} = -\nabla p_1(x, y, z, t) + \frac{1}{\mu_0} (\nabla \times \mathbf{B}_1(x, y, z, t)) \times \mathbf{B}_0 \\ \frac{\partial p_1(x, y, z, t)}{\partial t} = \gamma \frac{p_0}{\rho_0} \frac{\partial \rho_1(x, y, z, t)}{\partial t} \\ \frac{\partial \mathbf{B}_1(x, y, z, t)}{\partial t} = \nabla \times (\mathbf{v}_1(x, y, z, t) \times \mathbf{B}_0) \end{array} \right.$$

$$\mathbf{B}_0 = \mathbf{0}, \mathbf{B}_1 = \mathbf{0} \longrightarrow$$



$$\left\{ \begin{array}{l} \frac{\partial \rho_1(x, y, z, t)}{\partial t} = -\rho_0 \nabla \cdot \mathbf{v}_1(x, y, z, t) \\ \rho_0 \frac{\partial \mathbf{v}_1(x, y, z, t)}{\partial t} = -\nabla p_1(x, y, z, t) \\ \frac{\partial p_1(x, y, z, t)}{\partial t} = \gamma \frac{p_0}{\rho_0} \frac{\partial \rho_1(x, y, z, t)}{\partial t} \end{array} \right.$$

Dispersion relation of sound wave

$$\begin{aligned}\frac{\partial \rho_1(x, y, z, t)}{\partial t} &= -\rho_0 \nabla \cdot \mathbf{v}_1(x, y, z, t) \\ \rho_0 \frac{\partial \mathbf{v}_1(x, y, z, t)}{\partial t} &= -\nabla p_1(x, y, z, t) \\ \frac{\partial p_1(x, y, z, t)}{\partial t} &= \gamma \frac{p_0}{\rho_0} \frac{\partial \rho_1(x, y, z, t)}{\partial t}\end{aligned}$$

Partial differential equations

$$\begin{aligned}\rho_1(x, y, z, t) &= \rho^* e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} \\ \mathbf{v}_1(x, y, z, t) &= \mathbf{v}^* e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} \\ p_1(x, y, z, t) &= p^* e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)}\end{aligned}$$

$$\omega \rho_0 \mathbf{v}^* = k p^* \Rightarrow \mathbf{v}^* \parallel \mathbf{k} \dots \text{longitudinal wave}$$

apply $\mathbf{k} \cdot$ to both hands

Gas pressure gradient force is isotropic when p_0 is uniform, so we only consider the direction of $\mathbf{k}$.

$$-i \omega \rho^* = -\rho_0 i \mathbf{k} \cdot \mathbf{v}^*$$

$$-i \omega \rho_0 \mathbf{v}^* = -i \mathbf{k} p^*$$

$$-i \omega p^* = \gamma \frac{p_0}{\rho_0} (-i \omega \rho^*)$$

Algebraic equations

e^i -form

$$\frac{\partial}{\partial t} ?_1 \Rightarrow -i \omega ?_1$$

$$\nabla \cdot ?_1 \Rightarrow i \mathbf{k} \cdot ?_1$$

$$\nabla ?_1 \Rightarrow i \mathbf{k} ?_1$$

differential operation
 $\Rightarrow$ algebraic operation

$$-\omega \rho^* = -\rho_0 \mathbf{k} \cdot \mathbf{v}^*$$

$$\omega \rho_0 \mathbf{k} \cdot \mathbf{v}^* = k^2 p^*$$

$$p^* = \gamma \frac{p_0}{\rho_0} \rho^*$$

$$\begin{bmatrix} -\omega & \rho_0 & 0 \\ 0 & \omega \rho_0 - k^2 & \\ \gamma \frac{p_0}{\rho_0} & 0 & -1 \end{bmatrix} \begin{bmatrix} \rho^* \\ \mathbf{k} \cdot \mathbf{v}^* \\ p^* \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -\omega & \rho_0 & 0 \\ 0 & \omega \rho_0 - k^2 & \\ \gamma \frac{p_0}{\rho_0} & 0 & -1 \end{bmatrix} \begin{bmatrix} \rho^* \\ k \cdot v^* \\ p^* \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$k \cdot v^* \neq 0 \Rightarrow \nabla \cdot v_1 \neq 0 \dots$ *compressional wave*

if $k \cdot v^* = 0$, then

$$\rho^* = \frac{\rho_0}{\omega} k \cdot v^* = 0$$

$$p^* = \gamma \frac{p_0}{\rho_0} \rho^* = 0$$

requires $\det \begin{bmatrix} -\omega & \rho_0 & 0 \\ 0 & \omega \rho_0 - k^2 & \\ \gamma \frac{p_0}{\rho_0} & 0 & -1 \end{bmatrix} = 0$

$$\det \begin{bmatrix} -\omega & \rho_0 & 0 \\ 0 & \omega \rho_0 - k^2 & \\ \gamma \frac{p_0}{\rho_0} & 0 & -1 \end{bmatrix} = -\omega \omega \rho_0 (-1) + (\rho_0) (-k^2) \left(\gamma \frac{p_0}{\rho_0} \right) - 0 = \omega^2 \rho_0 - \gamma \frac{p_0}{\rho_0} \rho_0 k^2 = 0$$

$$\therefore \omega^2 - k^2 c_{s0}^2 = 0, \quad c_{s0} = \sqrt{\gamma \frac{p_0}{\rho_0}} \text{ (adiabatic sound speed)}$$