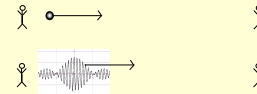


What is a wave?

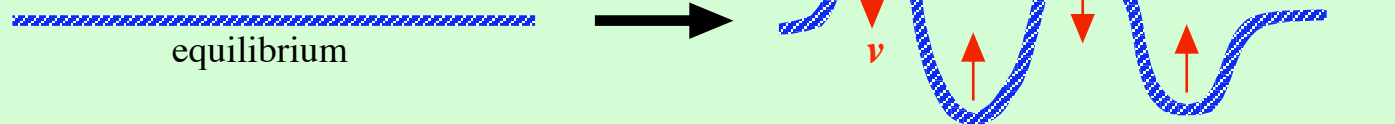
It is a **time-dependent phenomenon** in which a **physical state** is transported from one location to other locations through a **medium** (what is transported is **physical state**, not **physical object**).



Example:

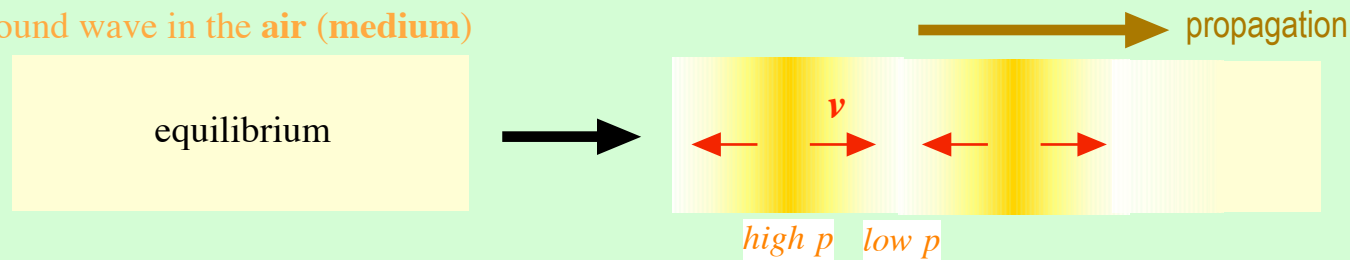
Transverse wave... oscillation direction $\perp$ propagation direction

Wave along a **string (medium)**



Longitudinal wave... oscillation direction $\parallel$ propagation direction

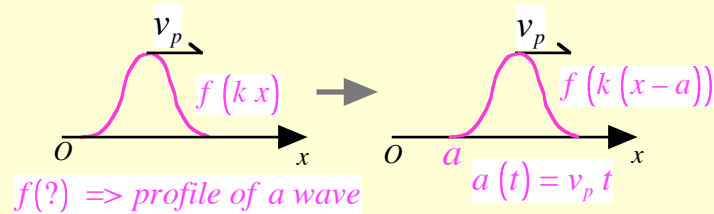
Sound wave in the **air (medium)**



Mathematical description of a wave

1-dimensional case

$$f(k(x - a(t))) = f(k(x - \boxed{v_p} t)) = f(kx - \omega t)$$



v_p : phase speed

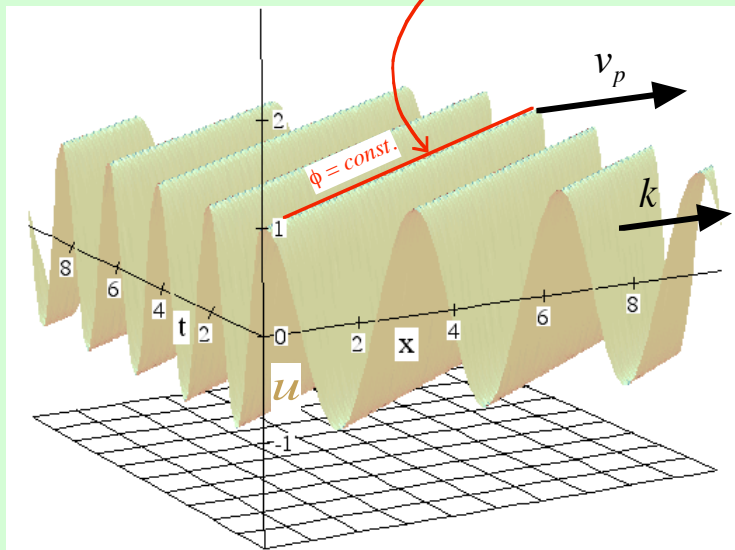
k : wavenumber ($\lambda = \frac{2\pi}{k}$: wavelength)

$\omega = k v_p$: angular frequency ($\mathcal{T} = \frac{2\pi}{\omega}$: period)

$\phi \equiv kx - \omega t$: phase $\Rightarrow$ physical state

Example: $u(x, t) = f(kx - \omega t) = \cos(2x - 3t)$

Trajectory of the same phase (= the same physical state) in (x, t) -space:
 $\phi \equiv kx - \omega t = \text{const.}$



$$\frac{d\phi}{dt} = 0 \text{ (Lagrangian derivative)} \Rightarrow k \frac{dx}{dt} - \omega = 0$$

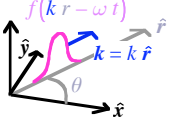
$$v_p = \frac{dx}{dt} = \frac{\omega}{k}: \text{phase speed}$$

Propagation speed of the same phase (the same physical state)

$$\cos(2x - 3t) = \cos\left(2\left[x - \boxed{\frac{3}{2}}t\right]\right)$$

The profile of " $\cos(2x)$ " moves toward positive x -direction at a speed of $3/2$.

3-dimensional case



$$\begin{aligned}
 k \cdot r &= k \hat{r} \cdot \mathbf{r} \\
 &= k (\cos \theta \hat{x} + \sin \theta \hat{y}) \cdot (x \hat{x} + y \hat{y}) \\
 &= k_x x + k_y y = \mathbf{k} \cdot \mathbf{r}
 \end{aligned}$$

$$u(x, y, z, t) = u^* f(\mathbf{k} \cdot \mathbf{r} - \omega t + \phi_0)$$

u^* : amplitude $\mathbf{r} = (x, y, z)$ $\mathbf{k} = (k_x, k_y, k_z)$: wavenumber vector $\Rightarrow$ propagation direction

ω : angular frequency $\mathbf{v}_p = \frac{\omega}{k} \hat{\mathbf{k}}$: phase velocity $\phi \equiv \mathbf{k} \cdot \mathbf{r} - \omega t + \phi_0$: phase

$$\mathbf{k} \cdot \mathbf{r} = \begin{pmatrix} k_x \\ k_y \\ k_z \end{pmatrix} \cdot \begin{pmatrix} x \\ y \\ z \end{pmatrix} = k_x x + k_y y + k_z z$$

$$0 = \frac{d\phi}{dt} = \frac{d}{dt} (\mathbf{k} \cdot \mathbf{r} - \omega t + \phi_0) \Rightarrow \mathbf{k} \cdot \frac{d\mathbf{r}}{dt} - \omega = 0 \Rightarrow \mathbf{k} \cdot \mathbf{v}_p - \omega = 0 \Rightarrow k v_p = \omega \text{ (because of } \mathbf{k} \parallel \mathbf{v}_p \text{)}$$

$$\mathbf{v}_p = \frac{\omega}{k} \hat{\mathbf{k}} = \frac{\omega}{\sqrt{k_x^2 + k_y^2 + k_z^2}} \begin{pmatrix} \frac{k_x}{\sqrt{k_x^2 + k_y^2 + k_z^2}} \\ \frac{k_y}{\sqrt{k_x^2 + k_y^2 + k_z^2}} \\ \frac{k_z}{\sqrt{k_x^2 + k_y^2 + k_z^2}} \end{pmatrix}$$

Linear wave

Linear wave... small-amplitude wave

Decomposition of a wave:



Equilibrium part (time-independent) + Perturbed part (time-dependent)

$$\rho_0(x, y, z) \quad \mathbf{v}_0(x, y, z) \equiv \mathbf{0}$$

$$p_0(x, y, z) \quad \mathbf{B}_0(x, y, z)$$

$$\rho_1(x, y, z, t) \quad \mathbf{v}_1(x, y, z, t)$$

$$p_1(x, y, z, t) \quad \mathbf{B}_1(x, y, z, t)$$

$$\begin{aligned} \rho(x, y, z, t) &\Rightarrow \rho_0(x, y, z) + \rho_1(x, y, z, t) \\ \mathbf{v}(x, y, z, t) &\Rightarrow \mathbf{0} + \mathbf{v}_1(x, y, z, t) \\ p(x, y, z, t) &\Rightarrow p_0(x, y, z) + p_1(x, y, z, t) \\ \mathbf{B}(x, y, z, t) &\Rightarrow \mathbf{B}_0(x, y, z) + \mathbf{B}_1(x, y, z, t) \end{aligned}$$

Perturbed part... much smaller than equilibrium part ($\square_0 \gg \square_1$)

=> product of perturbed parts is neglected

$$\square_1 \times \bigcirc_1 \dots \text{2nd-order} \doteq 0$$



Linear approximation

Linearized MHD equations for linear waves

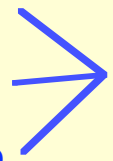
Assumption:

Equilibrium part

$$\rho_0(x, y, z) = \rho_0$$

$$p_0(x, y, z) = p_0$$

$$\mathbf{B}_0(x, y, z) = \mathbf{B}_0$$



uniform...

Plasma and magnetic field are uniformly distributed.

Substitute the decomposed form into ideal MHD equations (no diffusion):

$$\begin{aligned} \rho_0 + \rho_1(x, y, z, t) \\ \mathbf{0} + \mathbf{v}_1(x, y, z, t) \\ p_0 + p_1(x, y, z, t) \\ \mathbf{B}_0 + \mathbf{B}_1(x, y, z, t) \end{aligned}$$



$$\left\{ \begin{aligned} \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) &= 0 \\ \rho \left(\frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} \right) &= -\nabla p + \frac{1}{\mu_0} (\nabla \times \mathbf{B}) \times \mathbf{B} \\ \frac{\partial}{\partial t} \left(\frac{p}{\gamma - 1} \right) + \nabla \cdot \left(\frac{p}{\gamma - 1} \mathbf{v} \right) &= -p \nabla \cdot \mathbf{v} \quad \dots \text{adiabatic case} \\ \frac{\partial \mathbf{B}}{\partial t} &= \nabla \times (\mathbf{v} \times \mathbf{B}) \quad \dots \text{diffusion-free case} \end{aligned} \right.$$

Confirm these are the same.

$$\frac{d}{dt} \left(\frac{p}{\rho^\gamma} \right) = 0$$

Entropy

$$s = c_v \ln \left(\frac{p}{\rho^\gamma} \right) + \text{constant}$$

is conserved in the adiabatic case.

$$1. \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

$$\begin{matrix} \rho_0 + \rho_1(x, y, z, t) \\ \mathbf{0} + \mathbf{v}_1(x, y, z, t) \end{matrix}$$

$$\longrightarrow \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

$$\frac{\partial (\rho_0 + \rho_1(x, y, z, t))}{\partial t} + \nabla \cdot [(\rho_0 + \rho_1(x, y, z, t)) (\mathbf{0} + \mathbf{v}_1(x, y, z, t))] = 0$$

$$\nabla \cdot (\rho_0 \mathbf{v}_1(x, y, z, t) + \rho_1(x, y, z, t) \mathbf{v}_1(x, y, z, t)) \approx \nabla \cdot (\rho_0 \mathbf{v}_1(x, y, z, t))$$

2nd-order

$$\square_1 \bigcirc_1 \doteq 0$$

$$\frac{\partial \rho_0}{\partial t} = 0 \text{ } (\rho_0 \text{ is constant})$$

$(\rho_0 \text{ is constant})$

$$\frac{\partial \rho_1(x, y, z, t)}{\partial t}$$

$$\rho_0 \nabla \cdot \mathbf{v}_1(x, y, z, t)$$

Linearized mass conservation equation

$$\frac{\partial \rho_1(x, y, z, t)}{\partial t} = -\rho_0 \nabla \cdot \mathbf{v}_1(x, y, z, t)$$

$$2. \rho \left(\frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} \right) = -\nabla p + \frac{1}{\mu_0} (\nabla \times \mathbf{B}) \times \mathbf{B}$$

$$\begin{aligned} &\rho_0 + \rho_1(x, y, z, t) \\ &\mathbf{0} + \mathbf{v}_1(x, y, z, t) \\ &p_0 + p_1(x, y, z, t) \\ &\mathbf{B}_0 + \mathbf{B}_1(x, y, z, t) \end{aligned}$$

$$\longrightarrow \rho \left(\frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} \right) = -\nabla p + \frac{1}{\mu_0} (\nabla \times \mathbf{B}) \times \mathbf{B}$$

Left-hand side:

$$\rho \left(\frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} \right) \longrightarrow (\rho_0 + \rho_1(x, y, z, t)) \left[\frac{\partial \mathbf{v}_1(x, y, z, t)}{\partial t} + (\mathbf{v}_1(x, y, z, t)) \cdot \nabla \mathbf{v}_1(x, y, z, t) \right]$$

2nd-order
 $\square_1 \bigcirc_1 \doteq 0$
2nd-order
 $\square_1 \bigcirc_1 \doteq 0$

$$\longrightarrow \rho_0 \frac{\partial \mathbf{v}_1(x, y, z, t)}{\partial t}$$

Right-hand side:

$$-\nabla p + \frac{1}{\mu} (\nabla \times \mathbf{B}) \times \mathbf{B}$$

$$\rightarrow -\nabla(p_0 + p_1(x, y, z, t)) + \frac{1}{\mu_0} [\nabla \times (\mathbf{B}_0 + \mathbf{B}_1(x, y, z, t))] \times (\mathbf{B}_0 + \mathbf{B}_1(x, y, z, t))$$

$(p_0 \text{ is constant}) \quad (\mathbf{B}_0 \text{ is constant})$

$\underbrace{\hspace{10em}}_{\text{2nd-order}} \quad \square_1 \bigcirc_1 \doteq 0$

$$\rightarrow -\nabla p_1(x, y, z, t) + \frac{1}{\mu_0} (\nabla \times \mathbf{B}_1(x, y, z, t)) \times \mathbf{B}_0$$

Left-hand side

Right-hand side

$$\rho_0 \frac{\partial \mathbf{v}_1(x, y, z, t)}{\partial t} = -\nabla p_1(x, y, z, t) + \frac{1}{\mu_0} (\nabla \times \mathbf{B}_1(x, y, z, t)) \times \mathbf{B}_0$$

Linearized momentum equation