

Roles of Electric Field & Magnetic Field in MHD

Role of Electric Field

It drives evolution of magnetic field.

$$\frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E}^{MHD}$$

Electric field in MHD => two major components

$$\frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E}^{MHD} = \underbrace{-\nabla \times (-\mathbf{v} \times \mathbf{B})}_{\text{convective term}} - \underbrace{\nabla \times (\eta_{diff} \nabla \times \mathbf{B})}_{\text{diffusion term}}$$
$$\mathbf{E}_{conv} \equiv -\mathbf{v} \times \mathbf{B} \qquad \mathbf{E}_{resis} \equiv \eta_{diff} \nabla \times \mathbf{B}$$

- **Convective term:** $-\nabla \times \mathbf{E}_{conv} = \nabla \times (\mathbf{v} \times \mathbf{B})$

it drives evolution of magnetic field **via** macroscale motion (flow) of fluid elements.

- **Diffusion term:** $-\nabla \times \mathbf{E}_{resis} = -\nabla \times (\eta_{diff} \nabla \times \mathbf{B})$

it drives evolution of magnetic field **via** microscale motion (collision) of particles.

Convective term: $-\nabla \times \boldsymbol{E}_{conv} = \nabla \times (\boldsymbol{v} \times \boldsymbol{B})$

Convective limit...

$$\frac{\partial \mathbf{B}'}{\partial t'} = \nabla' \times \left(\mathbf{v}' \times \mathbf{B}' - \frac{1}{Rm} \nabla' \times \mathbf{B}' \right) \quad \dots \text{dimensionless form}$$

$$Rm \equiv \frac{v_0 l_0}{\eta_{diff}} \gg 1 \quad \longrightarrow \quad \text{we can neglect diffusion}$$

Three conditions for convective limit:

- I. typical length is very large
- II. typical velocity is very large
- III. magnetic diffusivity is small

e.g. solar corona...

condition I: $l_0 \sim 10^7 \text{ m}$

condition II: $v_0 \sim 10^5 \text{ m/s}$

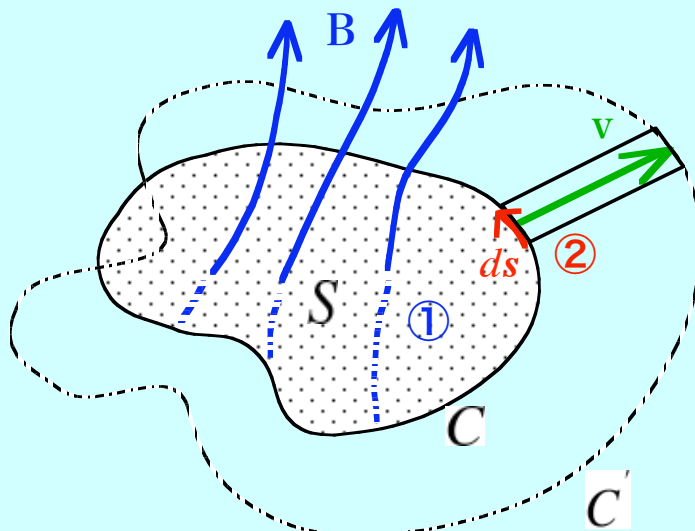
condition III: $\eta_{diff} = 5.2 \times 10^7 \frac{\ln N_D}{T^{3/2}} \sim 0.1 \text{ (m}^2/\text{s)}$

$$\longrightarrow Rm \sim 10^{13} \gg 1$$

Evolution of magnetic field via $\mathbf{E}_{conv} \equiv -\mathbf{v} \times \mathbf{B}$:

$$\Phi_S = \iint_S \mathbf{B} \cdot d\mathbf{S}$$

... total magnetic flux crossing S



Lagrangian derivative

Change of B when C is fixed ①
+
Change of C when B is fixed ②

$$\begin{aligned} \frac{d\Phi_S}{dt} &= \iint_S \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S} + \oint_C \mathbf{B} \cdot (\mathbf{v} \times d\mathbf{s}) \\ &= \iint_S \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S} + \oint_C (\mathbf{B} \times \mathbf{v}) \cdot d\mathbf{s} \\ &= \iint_S \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S} + \iint_S \nabla \times (\mathbf{B} \times \mathbf{v}) \cdot d\mathbf{S} \\ &= \iint_S \left[\frac{\partial \mathbf{B}}{\partial t} + \nabla \times (\mathbf{B} \times \mathbf{v}) \right] \cdot d\mathbf{S} \\ &= \iint_S \left[\frac{\partial \mathbf{B}}{\partial t} - \nabla \times (\mathbf{v} \times \mathbf{B}) \right] \cdot d\mathbf{S} \\ &= 0 \end{aligned}$$

scalar triple
vector product

Stokes's theorem

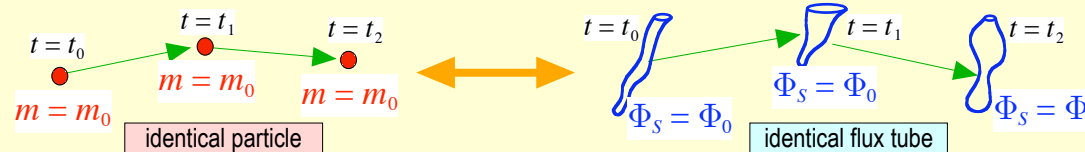
$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B}) \Rightarrow \frac{d\Phi_S}{dt} = 0$$

Physical meaning of $\frac{d\Phi_S}{dt} = 0 \Rightarrow$ Frozen-in evolution

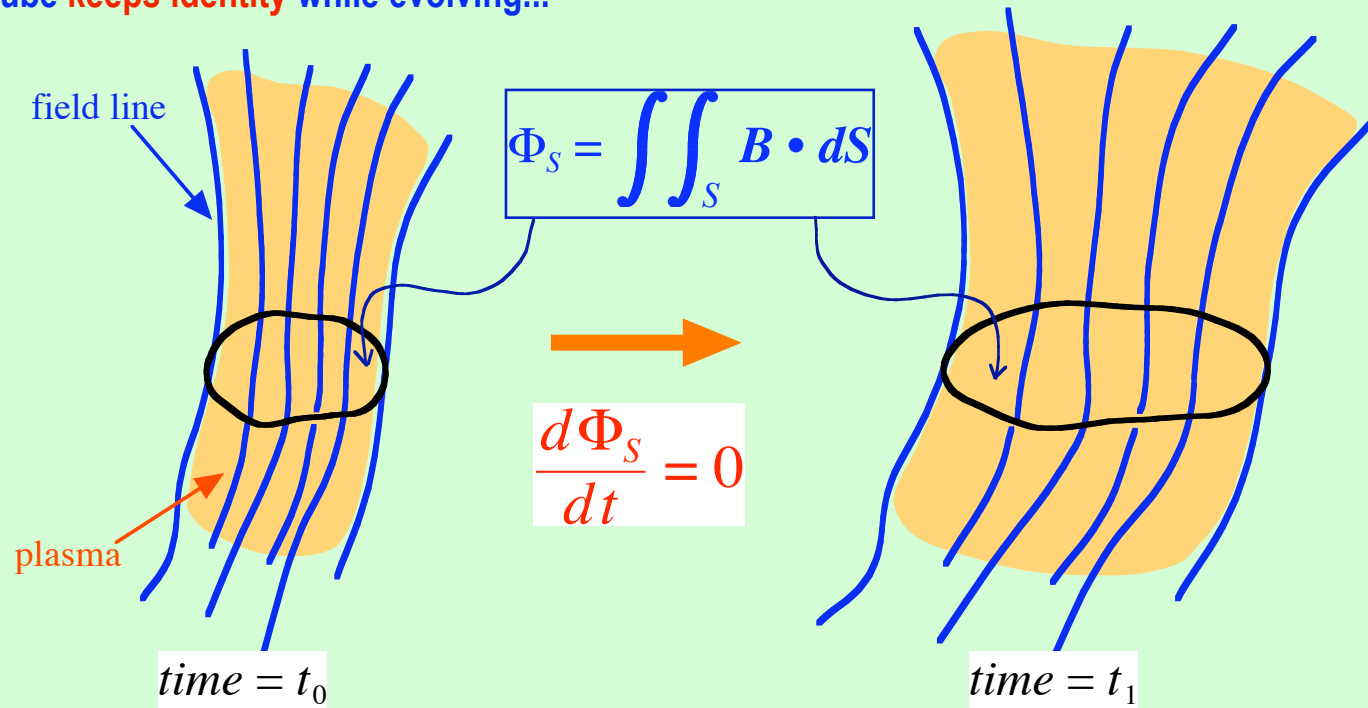
The number of magnetic field lines passing through the area bounded by any closed curve that moves together with a plasma does not change with time.



We can determine an identical flux tube in which magnetic field lines are frozen into the plasma.



Flux tube keeps identity while evolving...



Diffusion term: $-\nabla \times \boldsymbol{E}_{resis} = -\nabla \times (\eta_{diff} \nabla \times \boldsymbol{B})$

Diffusive limit...

$$\frac{\partial \mathbf{B}'}{\partial t'} = \nabla' \times \left(\mathbf{v}' \times \mathbf{B}' - \frac{1}{Rm} \nabla' \times \mathbf{B}' \right)$$

$$Rm \equiv \frac{v_0 l_0}{\eta_{diff}} \lesssim 1 \longrightarrow \text{we cannot neglect diffusion}$$

Three conditions for diffusive limit:

- I. typical length is very small
- II. typical velocity is very small
- III. magnetic diffusivity is large

In astronomical plasmas, their typical length l_0 is usually very large, so $Rm \equiv \frac{v_0 l_0}{\eta_{diff}} \lesssim 1$ is rare except for a region where magnetic field changes sharply (e.g. current sheet), which significantly reduces l_0 . It may also enhance η_{diff} via kinetic processes (e.g. wave-particle interaction).

Evolution of magnetic field via $E_{resis} \equiv \eta_{diff} \nabla \times B$:

When η_{diff} is constant, $\frac{\partial B}{\partial t} = -\nabla \times (\eta_{diff} \nabla \times B) \Rightarrow \frac{\partial B}{\partial t} = \eta_{diff} \nabla^2 B$

$$\frac{\partial q}{\partial t} = \eta_{diff} \nabla^2 q \quad \dots \text{diffusion equation}$$

$q = T, j_x, B_y, \dots$

Time scale of diffusion: τ_{diff}

$$\frac{\partial q}{\partial t} = \eta_{diff} \nabla^2 q \Rightarrow \frac{q}{\tau_{diff}} \sim \eta_{diff} \frac{q}{l_{diff}^2} \rightarrow \tau_{diff} \sim \frac{l_{diff}^2}{\eta_{diff}}$$

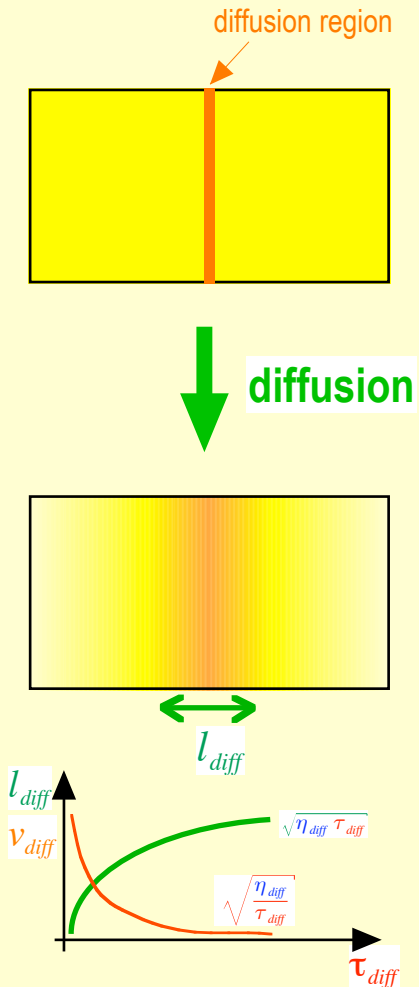
l_{diff} : width of diffusion region

Length scale of diffusion: $l_{diff} \sim \sqrt{\eta_{diff} \tau_{diff}}$

Velocity scale of diffusion: v_{diff}

$$v_{diff} \sim \frac{l_{diff}}{\tau_{diff}} = \frac{\sqrt{\eta_{diff} \tau_{diff}}}{\tau_{diff}} = \sqrt{\frac{\eta_{diff}}{\tau_{diff}}}$$

Diffusion proceeds fast at the beginning and then slowly.



Diffusion of an antiparallel magnetic field (*annihilation*)

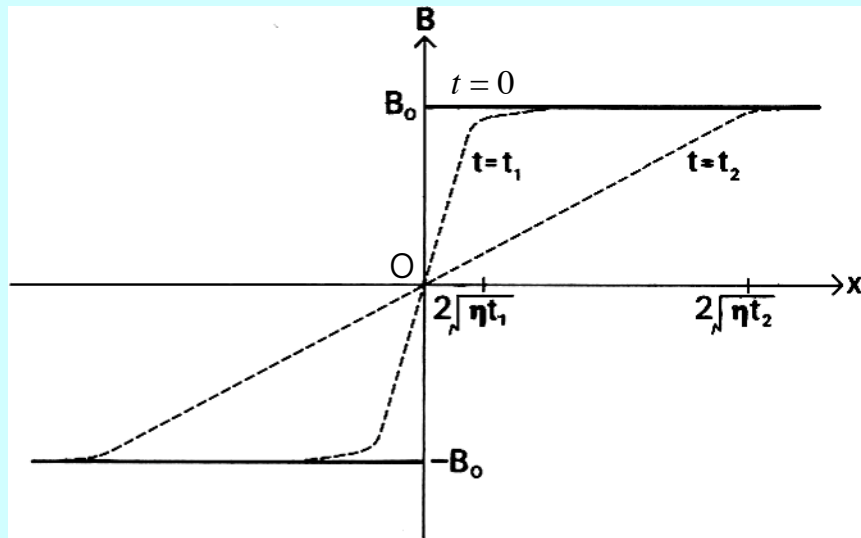
$$\frac{\partial B_y}{\partial t} = \eta_{diff} \frac{\partial^2 B_y}{\partial x^2}$$

$$B_y(x, t=0) = \begin{cases} B_0 & \text{for } x > 0 \\ -B_0 & \text{for } x < 0 \end{cases} \dots \text{initial condition}$$

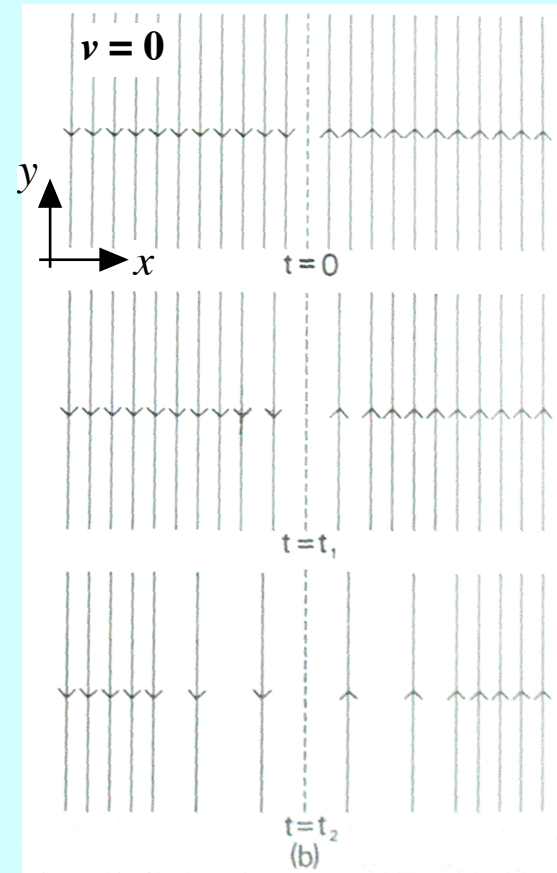
$$B_y(x = \pm \infty, t) = \begin{cases} B_0 & \text{for } x = \infty \\ -B_0 & \text{for } x = -\infty \end{cases} \dots \text{boundary condition}$$



$$B_y(x, t) = \frac{2 B_0}{\sqrt{\pi}} \operatorname{erf} \left(\frac{x}{\sqrt{4 \eta_{diff} t}} \right), \operatorname{erf}(\xi) = \int_0^\xi e^{-u^2} du$$



- magnetic field is annihilated at $x = 0$
- magnetic field diffuses through a plasma => **violates frozen-in evolution**



Solar MHD (Priest 1982)