

4) Energy equation

proton's energy conservation eq.
$$\frac{\partial}{\partial t} \left(\frac{n M}{2} u^{p2} + n M \mathcal{E}^p \right) + \frac{\partial}{\partial x_k} \left[\left(\frac{n M}{2} u^{p2} + n M \mathcal{E}^p + P^p \right) u_k^p - \pi_{ik}^p u_i^p + \mathcal{J}_k^p \right] - n M \left\langle \frac{F_k^p}{M} v_k^p \right\rangle = \int_v \frac{1}{2} M v^{p2} \left(\frac{\delta f^p}{\delta t} \right)_c d^3 v^p$$

electron's energy conservation eq.
$$\frac{\partial}{\partial t} \left(\frac{n m}{2} u^{e2} + n m \mathcal{E}^e \right) + \frac{\partial}{\partial x_k} \left[\left(\frac{n m}{2} u^{e2} + n m \mathcal{E}^e + P^e \right) u_k^e - \pi_{ik}^e u_i^e + \mathcal{J}_k^e \right] - n m \left\langle \frac{F_k^e}{m} v_k^e \right\rangle = \int_v \frac{1}{2} m v^{e2} \left(\frac{\delta f^e}{\delta t} \right)_c d^3 v^e$$

Add both equations (under the 0-th order approximation of $\frac{m}{M} \ll 1$)...

$$\left(\frac{n M}{2} u^{p2} + n M \mathcal{E}^p \right) + \left(\frac{n m}{2} u^{e2} + n m \mathcal{E}^e \right) \approx \frac{1}{2} \rho^{MHD} u^{MHD2} + \rho^{MHD} \mathcal{E}^{MHD}$$

$$\left(\frac{n M}{2} u^{p2} u_k^p \right) + \left(\frac{n m}{2} u^{e2} u_k^e \right) \approx \frac{1}{2} \rho^{MHD} u^{MHD2} u_k^{MHD}$$

assumption: $T^p \sim T^e \sim T^{MHD}$

$$\left[(n M \mathcal{E}^p + P^p) u_k^p \right] + \left[(n m \mathcal{E}^e + P^e) u_k^e \right] = \frac{\gamma}{\gamma-1} P^p u_k^p + \frac{\gamma}{\gamma-1} P^e u_k^e = \frac{\gamma}{\gamma-1} \frac{P^{MHD}}{2} [2 u_k^p - (u_k^p - u_k^e)] = \frac{\gamma}{\gamma-1} P^{MHD} \left(u_k^p - \frac{u_k^p - u_k^e}{2} \right)$$

$$(\pi_{ik}^p u_i^p) + (\pi_{ik}^e u_i^e) \approx \pi_{ik}^p u_i^p \approx \pi_{ik}^{MHD} u_i^{MHD}$$

assumption: $D_{ik}^p \sim D_{ik}^{MHD}$, $\mu^{MHD} \sim \mu^p \gg \mu^e (\sqrt{M} \gg \sqrt{m})$

$$(\mathcal{J}_k^p) + (\mathcal{J}_k^e) \approx \mathcal{J}_k^e \approx \mathcal{J}_k^{MHD}$$

assumption: $T^p \sim T^e \sim T^{MHD}$ and $\kappa^{MHD} \sim \kappa^e \gg \kappa^p (1/\sqrt{m} \gg 1/\sqrt{M})$

$$\left(n M \left\langle \frac{F_k^p}{M} v_k^p \right\rangle \right) + \left(n m \left\langle \frac{F_k^e}{m} v_k^e \right\rangle \right) = n e E_k (u_k^p - u_k^e) + n g_k (M u_k^p + m u_k^e) \approx E_k j_k^{MHD} + \rho^{MHD} g_k u_k^{MHD}$$

$$\approx (\rho^{MHD} \mathcal{E}^{MHD} + P^{MHD}) u_k^{MHD}$$

assumption:

$$\left| \frac{\partial}{\partial x_k} (P^{MHD} u_k^p) \right| \gg \left| \frac{k_B}{e} j_k \frac{\partial T}{\partial x_k} \right|$$

$$\int_v \frac{1}{2} M v^{p2} \left(\frac{\delta f^p}{\delta t} \right)_c d^3 v^p + \int_v \frac{1}{2} m v^{e2} \left(\frac{\delta f^e}{\delta t} \right)_c d^3 v^e = 0$$

Total energy is conserved via collisions between protons and electrons.

$$\rightarrow \frac{\partial}{\partial t} \left(\frac{\rho^{MHD}}{2} u^{MHD2} + \rho^{MHD} \mathcal{E}^{MHD} \right) + \frac{\partial}{\partial x_k} \left[\left(\frac{\rho^{MHD}}{2} u^{MHD2} + \rho^{MHD} \mathcal{E}^{MHD} + P^{MHD} \right) u_k^{MHD} - \pi_{ik}^{MHD} u_i^{MHD} + \mathcal{J}_k^{MHD} \right] = \mathbf{E}^{MHD} \cdot \mathbf{j}^{MHD} + \rho^{MHD} \mathbf{g} \cdot \mathbf{u}^{MHD}$$

Plasma energy equation

$\mathbf{E} \cdot \mathbf{j}$... represents electromagnetic contribution to the energy equation:

Electromagnetic energy equation:
$$\frac{\partial}{\partial t} \left(\frac{\epsilon_0 E^2}{2} + \frac{B^2}{2\mu_0} \right) + \nabla \cdot \left(\mathbf{E} \times \frac{\mathbf{B}}{\mu_0} \right) = - \mathbf{E} \cdot \mathbf{j}$$

Poynting flux source term
(plasma => electromagnetic field)

$$\mathbf{E}^{MHD} \cdot \mathbf{j}^{MHD} = - (\mathbf{u}^{MHD} \times \mathbf{B}) \cdot \mathbf{j}^{MHD} + \eta \mathbf{j}^{MHD} \cdot \mathbf{j}^{MHD} + \frac{1}{en} [\mathbf{j}^{MHD} \times \mathbf{B}] \cdot \mathbf{j}^{MHD}$$

source term
(electromagnetic field => plasma) = 0

$$\mathbf{E}^{MHD} \cdot \mathbf{j}^{MHD} = \eta j^{MHD^2} + (\mathbf{j}^{MHD} \times \mathbf{B}) \cdot \mathbf{u}^{MHD}$$

ηj^{MHD^2} ... **Ohmic heating rate** which reduces electromagnetic energy and increases internal energy of plasma => **source term in internal energy equation**

$(\mathbf{j}^{MHD} \times \mathbf{B}) \cdot \mathbf{u}^{MHD}$... **work rate done by $\mathbf{j} \times \mathbf{B}$ force**, which reduces electromagnetic energy and increases flow energy of plasma => **source term in flow energy equation**

$$\frac{\partial}{\partial t} (\rho^{MHD} \epsilon^{MHD}) + \frac{\partial}{\partial x_k} (\rho^{MHD} \epsilon^{MHD} u_k^{MHD}) = - P^{MHD} \frac{\partial u_k^{MHD}}{\partial x_k} - \frac{\partial \mathcal{S}^{MHD}}{\partial x_k} + \Psi^{MHD} + \eta j_{MHD}^2$$

Internal energy eq. in MHD

5) Charge conservation equation

$$\begin{array}{l}
 \left[\begin{array}{l} \frac{e}{M} \times \frac{\partial \rho^p}{\partial t} + \nabla \cdot (\rho^p \mathbf{u}^p) = 0 \\ \frac{e}{m} \times \frac{\partial \rho^e}{\partial t} + \nabla \cdot (\rho^e \mathbf{u}^e) = 0 \end{array} \right. \quad \begin{array}{l} \dots \text{proton's mass conservation } \rho^p \equiv n_p M \\ \dots \text{electron's mass conservation } \rho^e \equiv n_e m \end{array} \\
 \text{subtraction} \rightarrow \frac{\partial \rho_c}{\partial t} + \nabla \cdot \mathbf{j} = 0 \quad \begin{array}{l} \rho_c \equiv e (n_p - n_e) \\ \mathbf{j} \equiv e (n_p \mathbf{u}^p - n_e \mathbf{u}^e) \end{array}
 \end{array}$$

Charge conservation equation

In MHD, $n_p \sim n_e \sim n$, so $\rho_c^{MHD} \sim 0$

$$\nabla \cdot \mathbf{j}^{MHD} \sim 0 \quad \mathbf{j}^{MHD} \equiv ne (\mathbf{u}^p - \mathbf{u}^e)$$

Charge conservation equation in MHD

6) Equation of state

$$P^p = n k_B T^p \quad \dots \text{proton's equation of state}$$

$$P^e = n k_B T^e \quad \dots \text{electron's equation of state}$$

summation

$$P^{MHD} = 2 n k_B T^{MHD} = 2 \frac{\rho^{MHD}}{M + m} k_B T^{MHD} = \rho^{MHD} \frac{k_B}{\frac{M + m}{2}} T^{MHD} = \rho^{MHD} \frac{k_B}{\bar{m}} T^{MHD}$$

$$T^p \sim T^e \sim T^{MHD}$$

$$n = \frac{\rho^{MHD}}{M + m}$$

$$\bar{m} \equiv \frac{M + m}{2}$$

average mass

$$P^{MHD} = \rho^{MHD} \frac{k_B}{\bar{m}} T^{MHD}$$

Equation of state in MHD

Two-fluid dynamics equations (electron + proton)

$$\begin{array}{l} \rho^e, u_i^e, P^e, \mathcal{E}^e, T^e, \mathcal{I}_i^e, \Psi^e \\ \rho^p, u_i^p, P^p, \mathcal{E}^p, T^p, \mathcal{I}_i^p, \Psi^p \end{array}$$

$$\frac{\partial \rho^e}{\partial t} + \frac{\partial}{\partial x_k} (\rho^e u_k^e) = 0$$

$$\frac{\partial \rho^p}{\partial t} + \frac{\partial}{\partial x_k} (\rho^p u_k^p) = 0$$

$$\frac{\partial}{\partial t} (\rho^e u_i^e) + \frac{\partial}{\partial x_k} (\rho^e u_i^e u_k^e + P^e \delta_{ik} - \pi_{ik}^e) - \frac{\rho^e}{m^e} \langle F_i \rangle = \int_V m^e v_i^e \left(\frac{\delta f^e}{\delta t} \right)_c d^3 v$$

$$\frac{\partial}{\partial t} (\rho^p u_i^p) + \frac{\partial}{\partial x_k} (\rho^p u_i^p u_k^p + P^p \delta_{ik} - \pi_{ik}^p) - \frac{\rho^p}{m^p} \langle F_i \rangle = \int_V m^p v_i^p \left(\frac{\delta f^p}{\delta t} \right)_c d^3 v$$

$$\frac{\partial}{\partial t} \left(\frac{n m}{2} u^2 + n m \mathcal{E} \right) + \frac{\partial}{\partial x_k} \left[\left(\frac{n m}{2} u^2 + n m \mathcal{E} + P^e \right) u_k^e - \pi_{ik}^e u_i^e + \mathcal{I}_i^e \right] - n m \left\langle \frac{F_k}{m} v_i^e \right\rangle = \int_V \frac{1}{2} m v^2 \left(\frac{\delta f^e}{\delta t} \right)_c d^3 v^e$$

$$\frac{\partial}{\partial t} \left(\frac{n M}{2} u^2 + n M \mathcal{E}^p \right) + \frac{\partial}{\partial x_k} \left[\left(\frac{n M}{2} u^2 + n M \mathcal{E}^p + P^p \right) u_k^p - \pi_{ik}^p u_i^p + \mathcal{I}_i^p \right] - n M \left\langle \frac{F_k}{M} v_i^p \right\rangle = \int_V \frac{1}{2} M v^2 \left(\frac{\delta f^p}{\delta t} \right)_c d^3 v^p$$

$$P^e = \frac{\rho^e}{m^e} k_B T^e$$

$$P^p = \frac{\rho^p}{m^p} k_B T^p$$

Single fluid dynamics equations (MHD)

$$\rho^{MHD}, u_i^{MHD}, P^{MHD}, \mathcal{E}^{MHD}, T^{MHD}, \mathcal{I}_i^{MHD}, \Psi^{MHD}, j_i^{MHD}$$

mass conservation

$$\frac{\partial \rho^{MHD}}{\partial t} + \frac{\partial}{\partial x_k} (\rho^{MHD} u_k^{MHD}) = 0$$

summation

subtraction

$$\nabla \cdot \mathbf{j}^{MHD} \sim 0$$

charge conservation

momentum equation

$$\rho^{MHD} \left(\frac{\partial}{\partial t} u_i^{MHD} + u_k^{MHD} \frac{\partial}{\partial x_k} u_i^{MHD} \right) = (\mathbf{j}^{MHD} \times \mathbf{B})_i - \frac{\partial}{\partial x_k} (P^{MHD} \delta_{ik} - \pi_{ik}^{MHD})$$

summation

subtraction

$$\mathbf{E}^{MHD} + \mathbf{u}^{MHD} \times \mathbf{B} = \eta \mathbf{j}^{MHD} + \frac{1}{e n} \mathbf{j}^{MHD} \times \mathbf{B}$$

Ohm's law

Ohm's law $\cdot \mathbf{j}$

summation

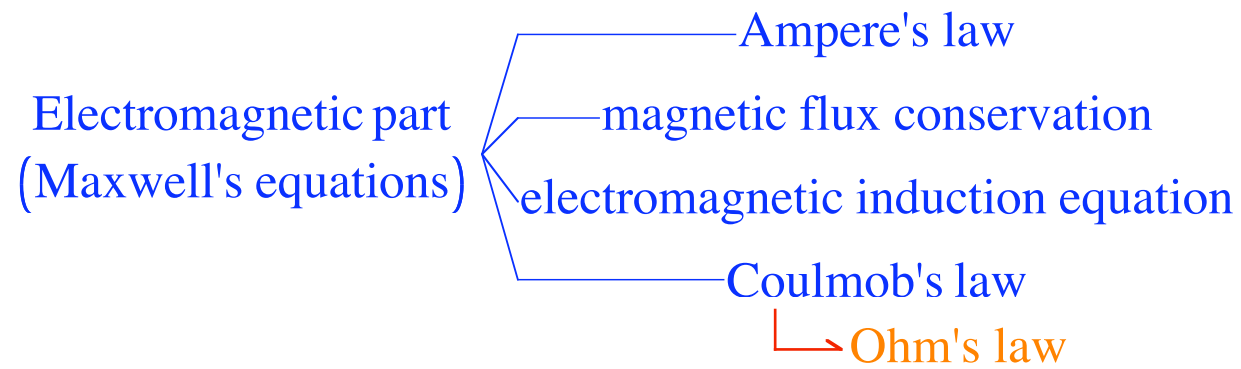
$$\frac{\partial}{\partial t} \left(\frac{\rho^{MHD}}{2} u^{MHD^2} + \rho^{MHD} \mathcal{E}^{MHD} \right) + \frac{\partial}{\partial x_k} \left[\left(\frac{\rho^{MHD}}{2} u^{MHD^2} + \rho^{MHD} \mathcal{E}^{MHD} + P^{MHD} \right) u_k^{MHD} - \pi_{ik}^{MHD} u_i^{MHD} + \mathcal{I}_i^{MHD} \right] = \mathbf{E}^{MHD} \cdot \mathbf{j}^{MHD} + \rho^{MHD} \mathbf{g} \cdot \mathbf{u}^{MHD}$$

energy equation

summation

$$P^{MHD} = \rho^{MHD} \frac{k_B}{\bar{m}} T^{MHD}$$

equation of state



(Fluid => Field)

Maxwell's equations (determine electromagnetic fields from the state of plasma)

MKS unit

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{j} + \frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t} \quad \dots \text{Generalized Ampere's law}$$

$$\nabla \cdot \mathbf{B} = 0 \quad \dots \text{Conservation of magnetic flux}$$

$$\nabla \times \mathbf{E} = - \frac{\partial \mathbf{B}}{\partial t} \quad \dots \text{Electromagnetic induction equation (Faraday's law)}$$

$$\nabla \cdot \mathbf{E} = \frac{\rho_c}{\epsilon_0} \quad \dots \text{Coulomb's law}$$

$$\epsilon_0 = 8.854 \times 10^{-12} \text{ F m}^{-1} \quad \dots \text{permittivity of free space}$$

$$\mu_0 = 4\pi \times 10^{-7} \text{ H m}^{-1} \quad \dots \text{magnetic permeability}$$

$$c = \frac{1}{\sqrt{\epsilon_0 \mu_0}} = 2.998 \times 10^8 \text{ m s}^{-1} \quad \dots \text{speed of light}$$

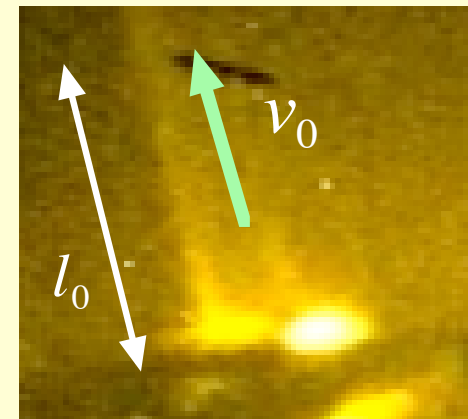
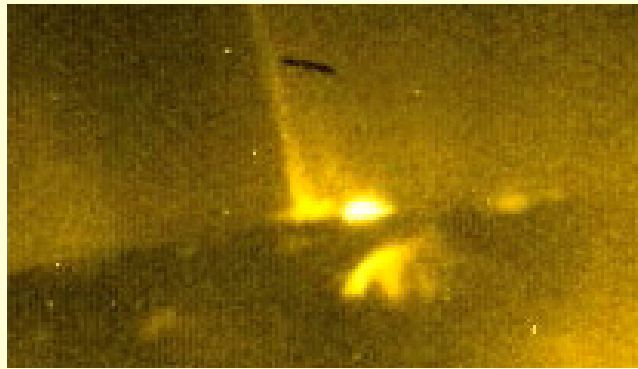
Units:

E ... V/m, B ... T (= 10^4 G), j ... A/m²

Typical scales in MHD

MHD phenomena... large-scale, slowly evolving phenomena

e.g. solar coronal jet



Typical scale:

length... $l_0 \sim 10^4$ km

velocity... $v_0 \sim 100$ km/s

time... $t_0 = \frac{l_0}{v_0} \sim 100$ s

r_B (gyration radius for a proton) ~ 1 m
 t_B (gyration period for a proton) $\sim 10^{-4}$ s

Approximations used in MHD

- $v_0 \ll c$ (non-relativistic approximation)

Comparison between electromagnetic terms in the generalized Ampere's law

$$\mathbf{j} = \frac{1}{\mu_0} \left(\nabla \times \mathbf{B} - \frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t} \right)$$

$$\frac{B_0}{l_0}$$

$$\frac{1}{c^2} \frac{E_0}{t_0}$$

Faraday's law

$$\nabla \times \mathbf{E} = - \frac{\partial \mathbf{B}}{\partial t} \rightarrow \frac{E_0}{l_0} \sim \frac{B_0}{t_0}$$

$$\frac{1}{c^2} \frac{E_0}{t_0} \approx \frac{1}{c^2} \frac{1}{t_0} \frac{l_0}{t_0} B_0 = \frac{1}{c^2} \frac{l_0^2}{t_0^2} \frac{B_0}{l_0} = \frac{v_0^2}{c^2} \frac{B_0}{l_0} \ll \frac{B_0}{l_0}$$

$$\therefore \frac{B_0}{l_0} \gg \frac{1}{c^2} \frac{E_0}{t_0} \Rightarrow \text{we neglect } \frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t}$$

Electromagnetic wave is NOT described by MHD.

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{j} + \frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t} \Rightarrow \nabla \times \mathbf{B} = \mu_0 \mathbf{j}$$

... **Ampere's law in MHD**

(determines magnetic field from current; consistent with MHD charge conservation eq.)